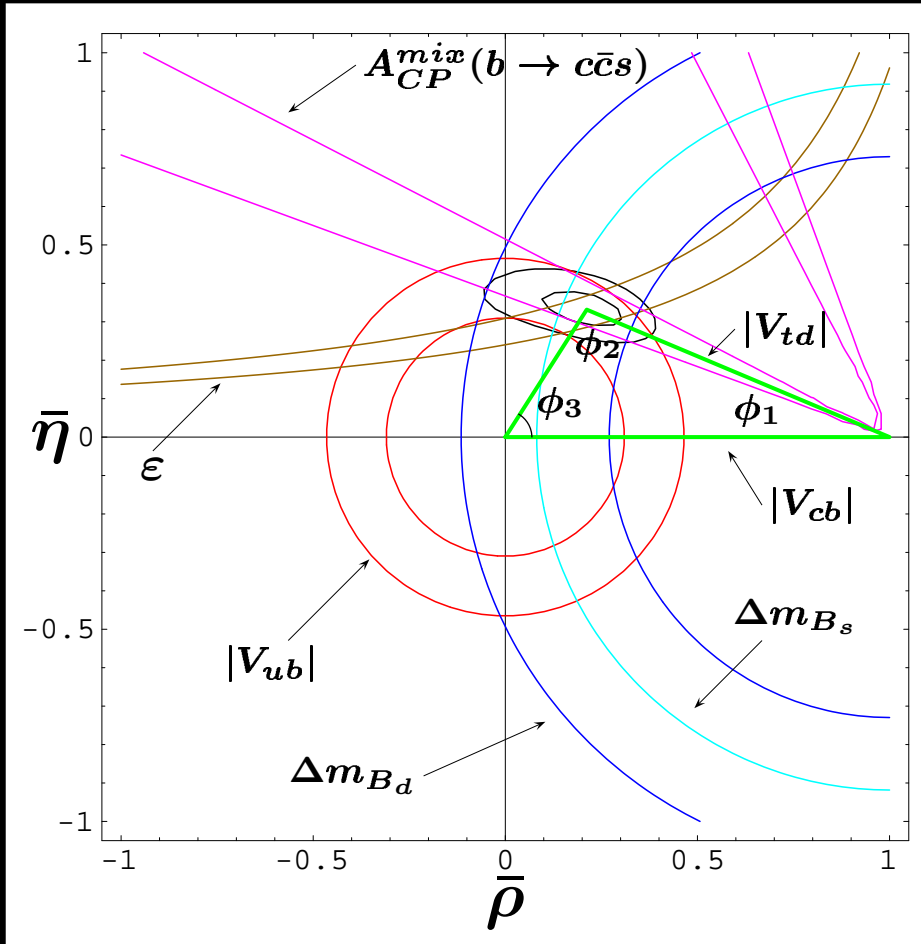


Physics with 10^9 and More B Mesons

Minoru TANAKA

Department of Physics
Osaka University

1. Introduction



- CPV in $K \rightarrow \pi\pi$
 ϵ parameter
- $b \rightarrow c, u$ transition
 $|V_{cb}|, |V_{ub}|$
- $B_d - \bar{B}_d$ mixing (Δm_{B_d})
 $|V_{td}|$
- $B_s - \bar{B}_s$ mixing (Δm_{B_s})
 $|V_{td}|$ (upper bound)
- CPA in $B_d \rightarrow \psi K_s, \dots$
($A_{CP}^{mix}(b \rightarrow c\bar{c}s)$)
 $\sin 2\phi_1$

The unitarity triangle looks closed.

But, this is not the whole story. We want to go further.

⇒ Precision test of the standard model

⇒ New physics search

Supersymmetry

The most attractive candidate of new physics

- Hierarchy problem: Why $M_{\text{weak}} \ll M_{\text{GUT}}$?
SUSY protects this hierarchy against radiative corrections
For this mechanism to work, $M_{\text{SUSY}} \sim M_{\text{weak}}$
- SUSY must be broken.
No degenerate SUSY pair of a fermion and a boson.
- SUSY breaking must be soft.
Cancellation of the second order divergence is required.
Allowed soft breakings: gaugino masses, scalar masses, trilinear scalar couplings, etc.
⇒ New sources of flavor and CP violation

Super B factory will unveil these new flavor and CP violation.

Super B factory

KEK

$$L \sim 10^{35} \text{ cm}^{-2} \text{ s}^{-1}, 1 \text{ ab}^{-1} / \text{yr}, 10^9 B\bar{B} \text{ pairs} / \text{yr}$$

SLAC

$$L \sim 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$$

PLAN

1. Introduction
2. Flavor structure of soft SUSY breakings
3. SUSY effects in $b \rightarrow s$ transitions
4. $\bar{B} \rightarrow D\tau\bar{\nu}$
5. Summary and discussion

2. Soft SUSY breakings

Scalar masses, scalar trilinear couplings

⇒ New sources of flavor and CP violation

- Squark mass matrix (6×6)

$$\tilde{M}_q^2 = \begin{pmatrix} \tilde{M}_{q,LL}^2 & \tilde{M}_{q,LR}^2 \\ \tilde{M}_{q,LR}^{2\dagger} & \tilde{M}_{q,RR}^2 \end{pmatrix}, \quad (q = u, d)$$

- Quark mass matrix (3×3)

$$V_L^q M_q V_R^{q\dagger} = \hat{M}_q \text{ (diagonal)}$$

FV is caused by off-diagonal

$$(\delta_M^{qN})_{ij} \equiv \frac{(V_M^q \tilde{M}_{q,MN}^2 V_N^{q\dagger})_{ij}}{\tilde{m}^2}$$

- $K-\bar{K}$ mixing

$$\sqrt{(\delta_{LL}^d)_{12}(\delta_{RR}^d)_{12}} < 0.006, \quad (\delta_{MM}^d)_{12} < 0.05$$

$$\sqrt{(\delta_{LL}^u)_{12}(\delta_{RR}^u)_{12}} < 0.04, \quad (\delta_{MM}^u)_{12} < 0.1$$

- $B-\bar{B}$ mixing

$$\sqrt{(\delta_{LL}^d)_{13}(\delta_{RR}^d)_{13}} < 0.04, \quad (\delta_{MM}^d)_{13} < 0.1$$

- $b \rightarrow s\gamma$

$$(\delta_{LR}^d)_{23} < 0.03 \quad (\tilde{m} \sim 400\text{GeV}) \quad \text{Kane et al., hep-ph/0212092}$$

$$\Rightarrow (\delta_{MN}^q)_{ij} < 0.01 \sim 0.1 \quad (\tilde{m} \sim 1\text{TeV})$$

We need some mechanism to suppress δ 's.

Note: no constraint on $(\delta_{LL,RR}^d)_{23}$

Possible mechanisms

Several ideas in the market

- Universal soft breakings
Gravity mediation of SUSY breaking (SUGRA),
Gauge mediation
Phenomenology depends on the scale of SUSY breaking and extra interaction below it (GUT, ν Yukawa).
- Flavor symmetry
Must explain the fermion masses and mixings.
Depends on the breaking sector.
- Alignment
- Effective SUSY

SUSY breaking mechanism $\iff \delta's \iff$ B physics

3. *SUSY* effects in $b \rightarrow s$ transitions

The large 2–3 mixing in the ν oscillation

- $B \rightarrow X_s \gamma$
 - $B \rightarrow X_s l^+ l^-$
 - $B \rightarrow X_s \nu \bar{\nu}$
 - $B \rightarrow \phi K_S$
 - $B \rightarrow \eta' K_S$
 - $B \rightarrow K \pi$
 - Δm_{B_s}
- etc.

$$B \rightarrow \phi K_S \quad (b \rightarrow s\bar{s}s)$$

- Experiment

$$\text{BR} = (8.4_{-2.1}^{+2.5}) \times 10^{-6} \quad (\text{BaBar} + \text{Belle})$$

$$S_{\phi K_S} = -0.39 \pm 0.41 \quad (\text{BaBar} + \text{Belle})$$

$$-A_{\phi K_S} = C_{\phi K_S} = 0.56 \pm 0.41 \pm 0.12 \quad (\text{Belle})$$

- SM prediction

$$\text{BR} \sim 5 \times 10^{-6}$$

$$S_{\phi K_S} = S_{\psi K_S} = 0.734 \pm 0.054 \quad (\text{theoretically clean})$$

$$C_{\phi K_S} \sim -0.008$$

- SUSY contribution to explain $S_{\phi K_S}$

larger hadronic uncertainty

$$|\text{Im}(\delta_{LR}^d)_{23}| \sim 2 \times 10^{-3} \quad (m_{\tilde{g}}/m_b \text{ enhancement})$$

$$|\text{Im}(\delta_{RL}^d)_{23}| \sim 2 \times 10^{-3} \quad (m_{\tilde{g}}/m_b \text{ enhancement})$$

$$|\text{Im}(\delta_{LL}^d)_{23}| \sim 0.4 \quad (\text{controversial})$$

$$|\text{Im}(\delta_{RR}^d)_{23}| \sim 0.4 \quad (\text{controversial})$$

\Rightarrow Correlation with other observables

$B \rightarrow \phi K_S$ (cont'd)

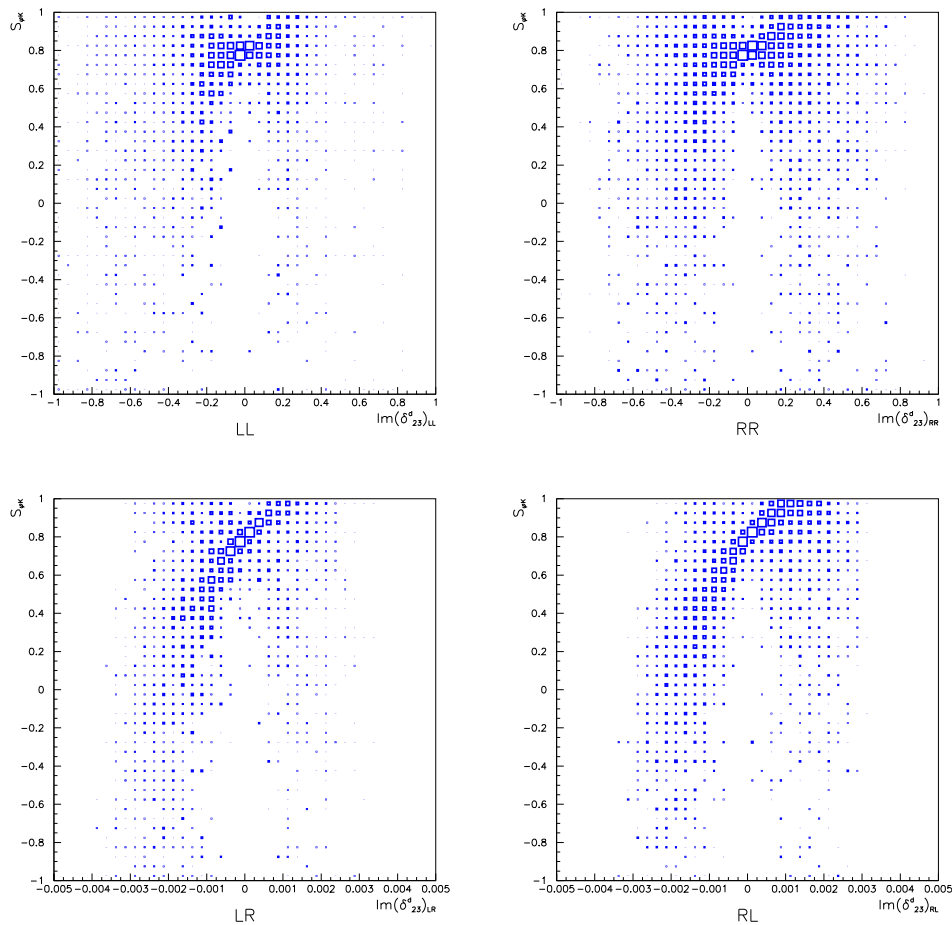


Figure 5: Correlations between $S_{\phi K}$ and $\text{Im}(\delta_{23}^d)_{AB}$ for $m_{\tilde{q}} = m_{\tilde{s}} = 350$ GeV and $AB = (L, RR, LR, RL)$. Constraints from $BR(B \rightarrow X_s \gamma)$, $A_{CP}(B \rightarrow X_s \gamma)$, $BR(B \rightarrow X_s l^+ l^-)$ and the lower bound on ΔM_s have been used.

$S_{\phi K_S}$

Ciuchini et al
hep-ph/0212397

$B \rightarrow X_s \gamma$

- Experiment

$$\text{BR} = (3.29 \pm 0.34) \times 10^{-4}, \quad A_{CP} = -0.02 \pm 0.04$$

- SM prediction

$$\text{BR} = (3.60 \pm 0.30) \times 10^{-4}$$

$$|A_{CP}| \lesssim 0.5\%, \quad |A_{CP}^{\text{mix}}| \sim 1\%$$

- SUSY contribution

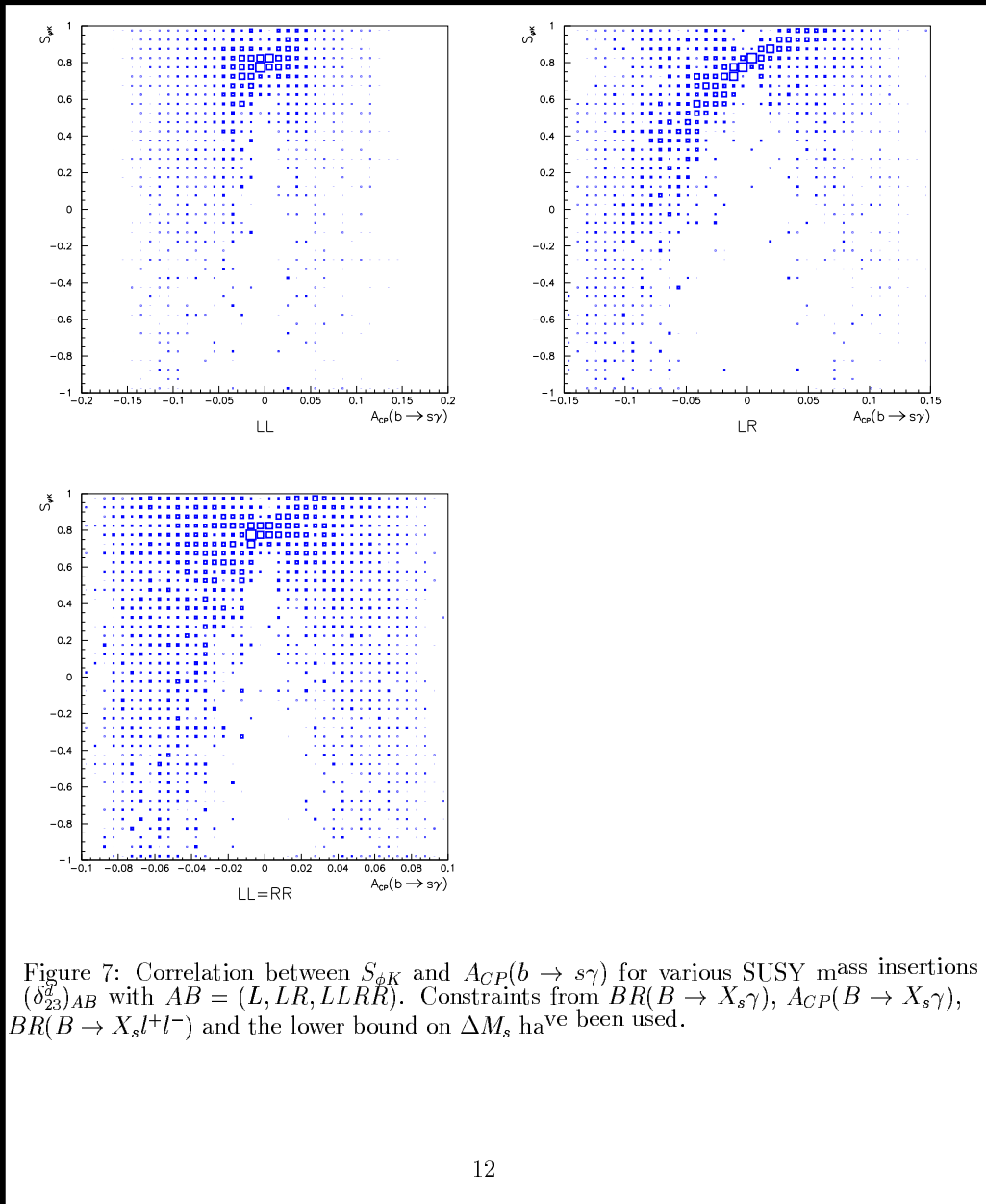
- ★ $|(\delta_{LR}^d)_{23}| \sim 0.002, b_R \rightarrow s_L$ as the SM
 \Rightarrow Interference with the SM part, $|A_{CP}| \lesssim 15\%$

- ★ $|(\delta_{RL}^d)_{23}| \sim 0.002, b_L \rightarrow s_R$
 $\Rightarrow |A_{CP}^{\text{mix}}| \sim O(0.1)$ or larger?

- ★ $|(\delta_{LL}^d)_{23}| \sim 0.4, b_R \rightarrow s_L$ as the SM
 $\Rightarrow |A_{CP}| \lesssim 15\%$

- ★ $|(\delta_{RR}^d)_{23}| \sim 0.4, b_L \rightarrow s_R$
 $\Rightarrow |A_{CP}^{\text{mix}}| \sim O(0.1)$ or larger?

$B \rightarrow \phi K_S$ and $B \rightarrow X_s \gamma$



$$A_{CP}(b \rightarrow s\gamma) - S_{\phi K_S}$$

Ciuchini et al
hep-ph/0212397

$$B \rightarrow \eta^{(\prime)} K_S, K\pi, \dots \quad (b \rightarrow s\bar{q}q)$$

- Experiment

$$S_{\eta' K_S} = 0.73 \pm 0.36_{0.06}^{+0.05} (\simeq S_{\psi K_S}) \text{ (Belle)}$$

$$A_{\eta' K_S} = 0.26 \pm 0.22 \pm 0.03 \text{ (Belle)}$$

- SUSY contribution

Similar as ϕK_S at the quark level, but different parity

$$\mathcal{P}(\bar{B}^0) = -1, \mathcal{P}(VP) = -1, \mathcal{P}(PP) = +1$$

$$\langle VP | \mathcal{O}_+ | \bar{B}^0 \rangle, \mathcal{O}_+ = (\bar{s}s)_V (\bar{s}b)_V, (\bar{s}s)_A (\bar{s}b)_A$$

$$\langle PP | \mathcal{O}_- | \bar{B}^0 \rangle, \mathcal{O}_- = (\bar{s}s)_A (\bar{s}b)_V, (\bar{s}s)_V (\bar{s}b)_A$$

$$VV(AA) \propto RR + LL \pm (LR + RL) \propto$$

$$(\delta_{RR}^d)_{23} + (\delta_{LL}^d)_{23}$$

$$VA(AV) \propto RR - LL \pm (LR - RL) \propto$$

$$(\delta_{RR}^d)_{23} - (\delta_{LL}^d)_{23}$$

\Rightarrow SUSY contributions to VP and PP are different.

Δm_{B_s}

- Experiment

$$\Delta m_{B_s} > 14.4 \text{ ps}^{-1} \text{ (95\% CL)}$$

- SM prediction

$$\Delta m_{B_s} = 17.3_{-0.7}^{+1.5} \text{ ps}^{-1}$$

- SUSY contribution

Could be larger (beyond Tevatron reach $\sim 41 \text{ ps}^{-1}$)

$$(\delta_{LL}^q) \sim 0.5 \Rightarrow \Delta m_{B_s} \sim 30 \text{ ps}^{-1}$$

Δm_{B_s} (cont'd)

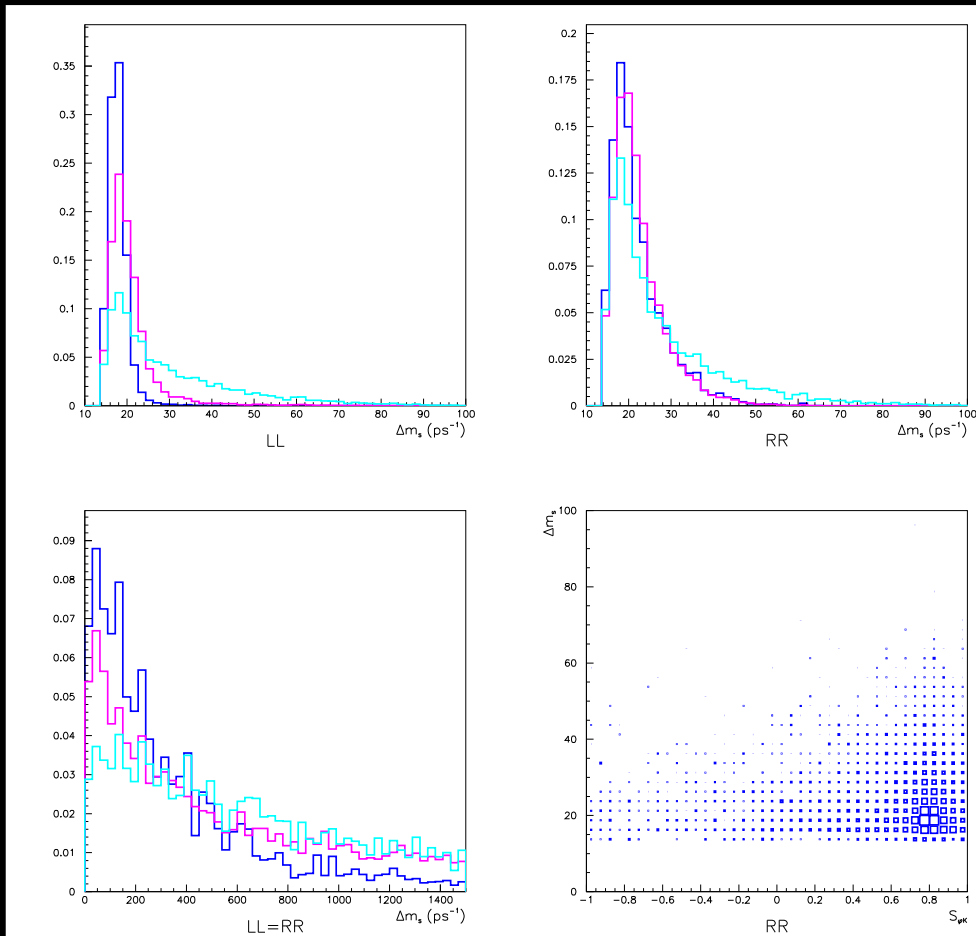


Figure 8: Distributions of ΔM_s for various SUSY mass insertions $(\delta_{23}^d)_{AB}$ with $AB = (L, RR, LLRR)$. Different curves correspond to the inclusion of constraints from $B \rightarrow X_s \gamma$ only (magenta), $B \rightarrow X_s l^+ l^-$ only (cyan) and all together (blue). Lower right: correlation between ΔM_s and $S_{\phi K}$ in the RR case.

Δm_{B_s}

Ciuchini et al
hep-ph/0212397

$$B \rightarrow X_s l^+ l^- \quad (b \rightarrow s l^+ l^-)$$

- Experiment

$$\text{BR} = (6.1 \pm 1.4_{-1.1}^{+1.3}) \times 10^{-6} \text{ (Belle)}$$

- SM prediction

Ali et al., PRD66,034002, 2002

$$\text{BR}(e^+ e^-) = (6.89 \pm 1.01) \times 10^{-6}$$

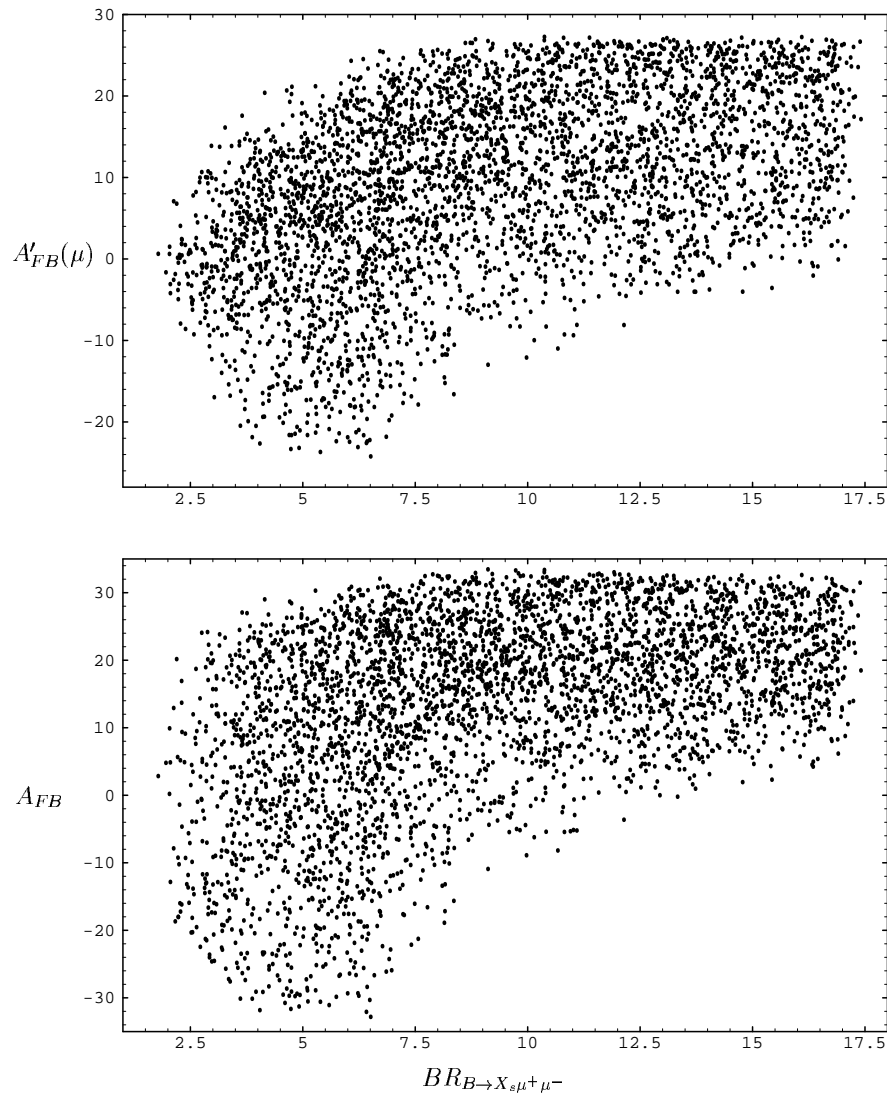
$$\text{BR}(\mu^+ \mu^-) = (4.15 \pm 0.70) \times 10^{-6}$$

$$A_{\text{FB}} \simeq 23\%, \quad A'_{\text{FB}}(\mu) \simeq 11\%$$

- SUSY contribution

$(\delta_{LL}^q)_{23}, (\delta_{LR}^q)_{23}$ are most relevant.

$B \rightarrow X_s l^+ l^-$ FB asymmetry



FB asym.

Lunghi, Scimemi
NP574, 43, 2000

Figure 1: Inclusive decay $B \rightarrow X_s \mu^+ \mu^-$. Scatter plots of the Integrated F-B asymmetries (A_{FB} and A'_{FB}) versus the BR (in units of 10^{-6}). Scenario A.

$B \rightarrow X_s l^+ l^-$ CP asymmetry

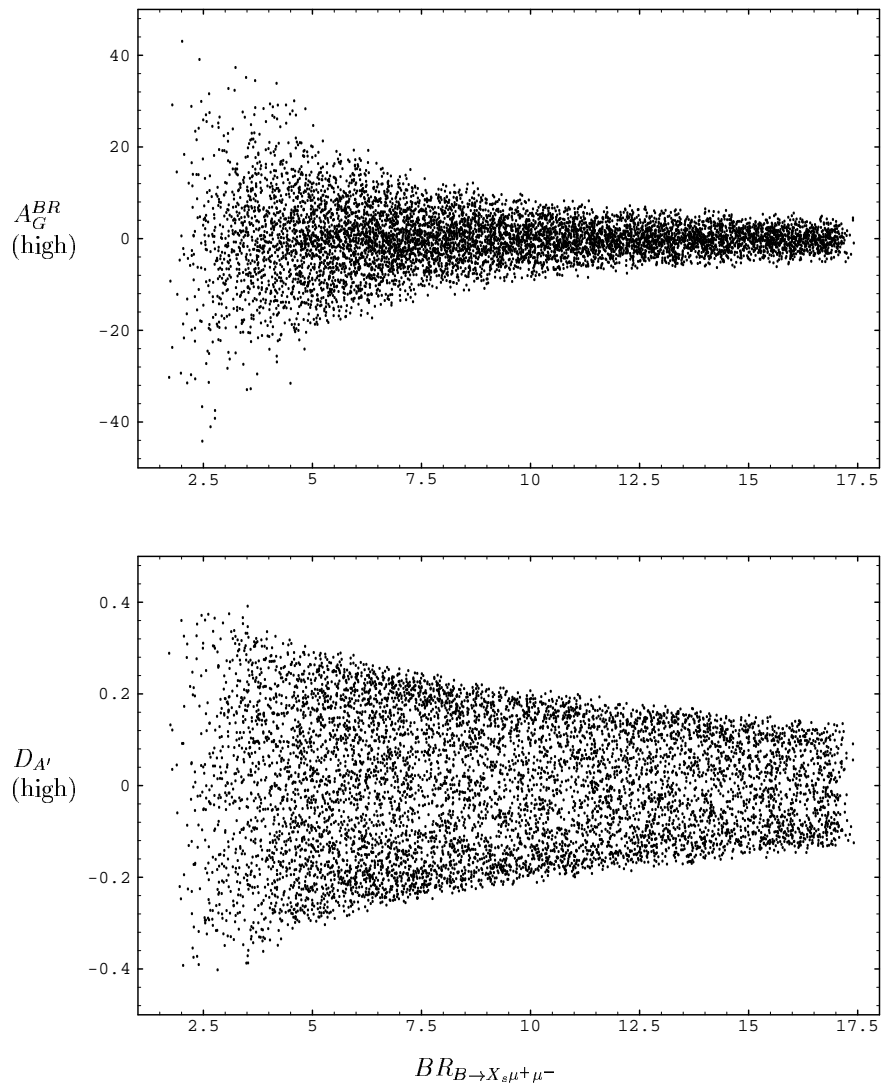


Figure 7: Inclusive decay $B \rightarrow X_s l^+ l^-$. Scatter plots of A_G^{BR} and $D_{A'}$ versus the integrated BR (in units of 10^{-6}) in the high- s region. Scenario A.

CP asym.
High s region
(14 – 23 GeV²)

Lunghi, Scimemi
NP574, 43, 2000

$$B \rightarrow X_s \nu \bar{\nu} \quad (b \rightarrow s \bar{\nu} \bar{\nu})$$

Less theoretical uncertainty.

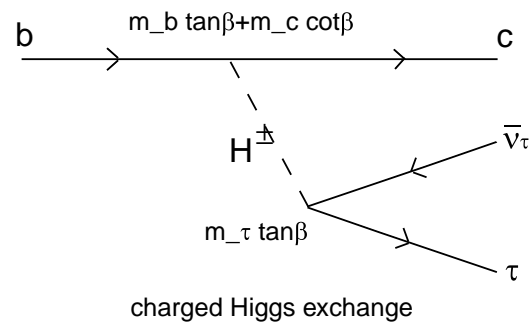
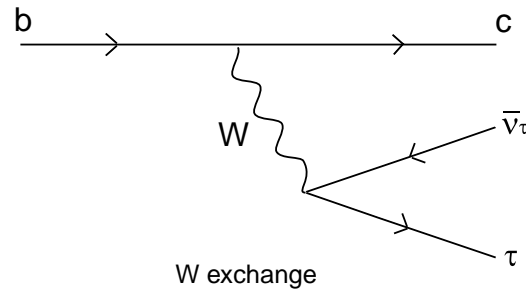
- Experiment
BR < 7.7×10^{-4} (ALEPH)
- SM prediction
BR = $(3.5 \pm 0.7) \times 10^{-5}$
- SUSY contribution
 $(\delta_{LL}^u)_{23}, (\delta_{LR}^u)_{23}$

4. $\bar{B} \rightarrow D\tau\bar{\nu}$

The best probe into the MSSM Higgs sector at B factories. Some recent theoretical results were reported by T. Miura in HLBF3.

3rd Workshop on Higher Luminosity B factory August 6-7, 2002

Feynman diagram of $\bar{B} \rightarrow D\tau\bar{\nu}$



August 6, 2002

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Why $\bar{B} \rightarrow D\tau\bar{\nu}_\tau$?

- Tree level process
Less dependent on the other sectors of the model
(Radiative correction is important when $\tan\beta \gg 1$)
- Charged Higgs couplings to fermions: $\propto m_f$

$$\begin{aligned} \text{Amp}(\bar{K} \rightarrow \pi\mu\bar{\nu}) &\sim m_s m_\mu \\ &\quad \Downarrow \sim 500 \text{ times} \\ \text{Amp}(\bar{B} \rightarrow D\tau\bar{\nu}) &\sim m_b m_\tau \end{aligned}$$

- B factories

$$\begin{array}{l} \# \text{ of B} \\ \sim 10^8 \end{array} \times \begin{array}{l} \text{BR}(b \rightarrow c\tau\bar{\nu}_\tau) \\ \sim 0.01 \end{array} \sim O(10^6) \text{ events !}$$

Super B factories \Rightarrow More events. BR, Distributions.

D vs D^*

$\bar{B} \rightarrow$	DW^*	DH^*	$D_L^*W^*$	$D_L^*H^*$
$L = 0$	H_s^s	H_s^s	H_0^0	
$L = 1$	H_0^s		H_s^0	H_s^0

$$H_0^s, H_s^0 \propto \sqrt{(m_B - m_M)^2 - q^2}$$

$$\longrightarrow 0 \quad \text{as } q^2 \rightarrow q_{max}^2 = (m_B - m_M)^2$$

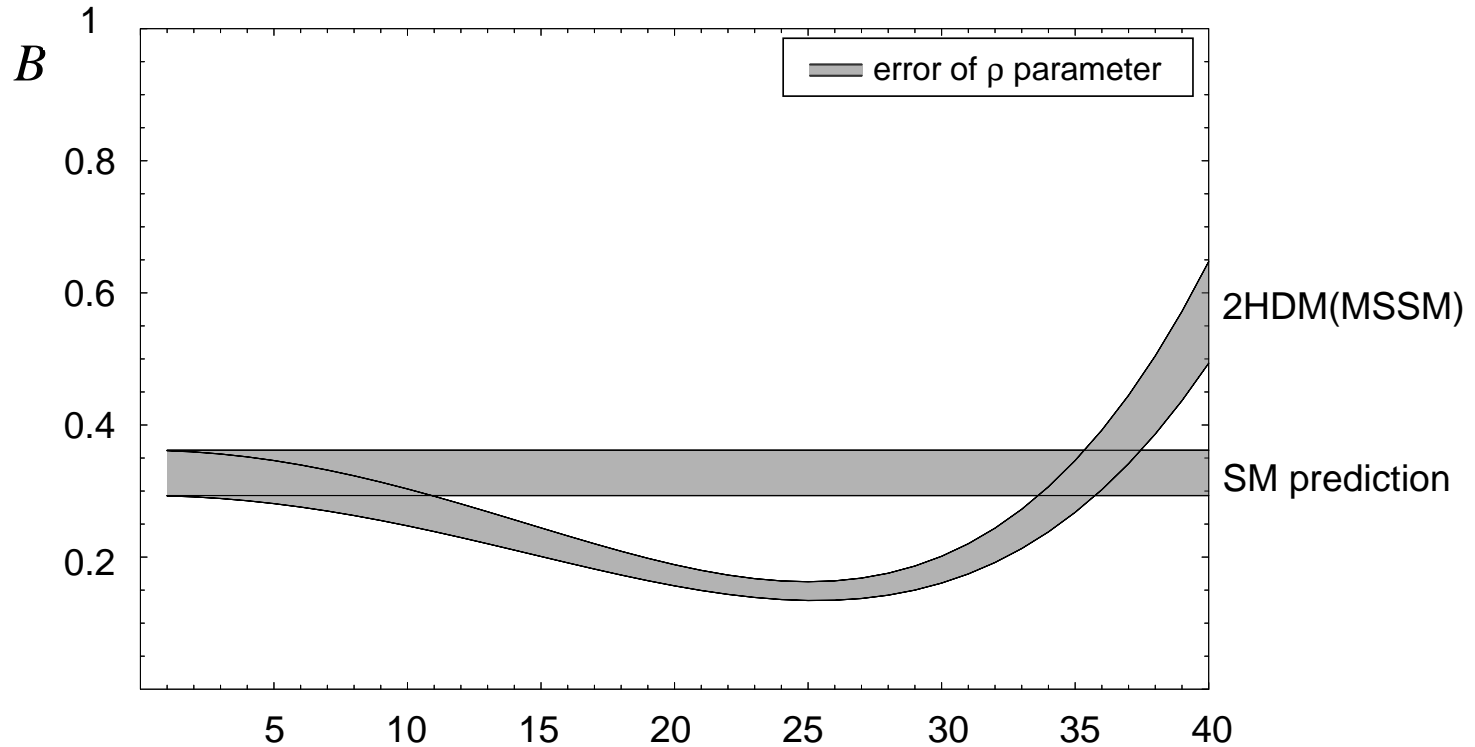
D is more sensitive than D_L^* .

Branching ratio

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$$\Lambda_{\overline{MS}} = 0.25(\text{GeV})$$

$$B(R) = \frac{\Gamma(\overline{B} \rightarrow D\tau\overline{\nu}_\tau)}{\Gamma(\overline{B} \rightarrow D\mu\overline{\nu}_\mu)_{SM}} \quad (\text{including error of } \rho \text{ parameter})$$



$$R = \frac{m_W}{m_H} \tan \beta$$

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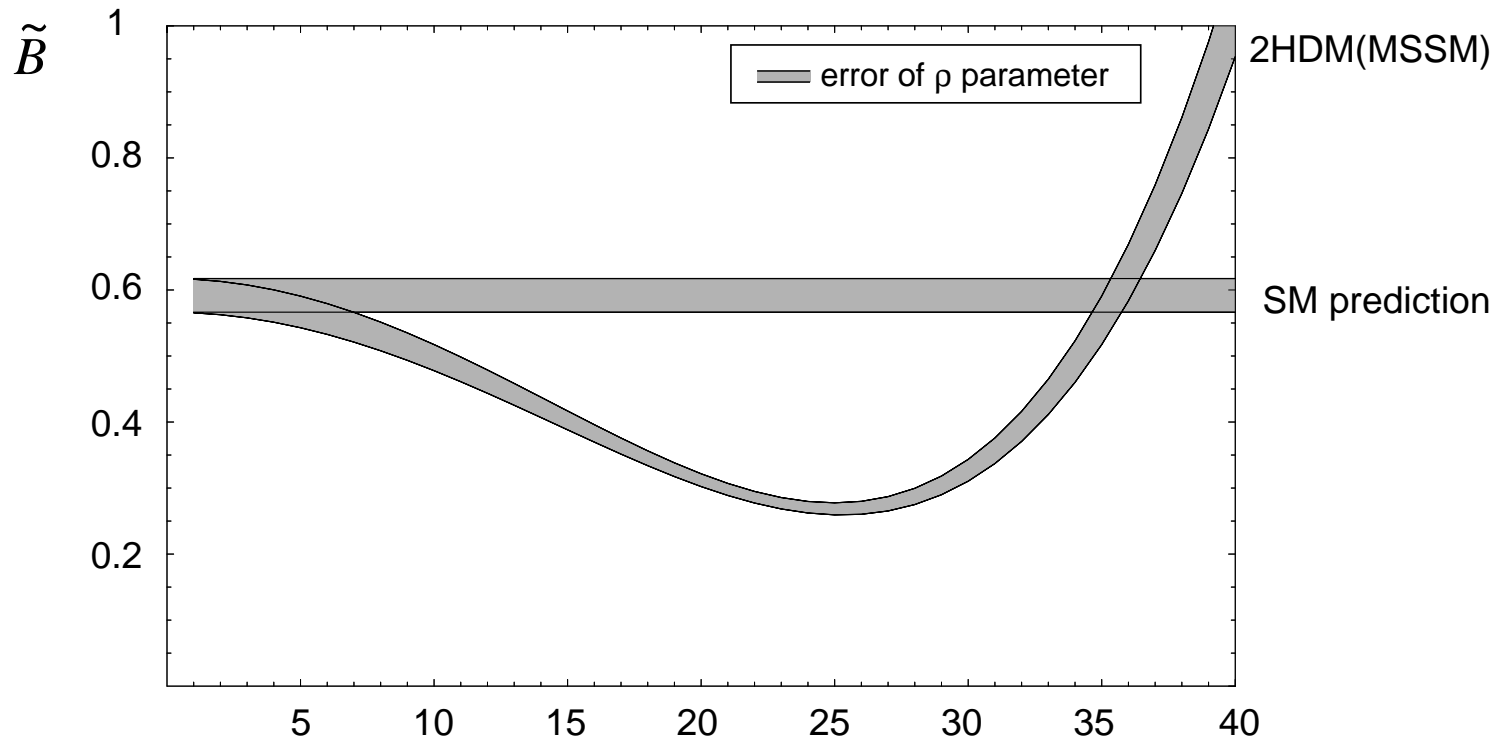
Branching ratio (cont'd)

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$$\Lambda_{\overline{MS}} = 0.25(\text{GeV})$$

$$\tilde{B}(R) = \frac{\Gamma(\overline{B} \rightarrow D\tau\bar{\nu}_\tau)}{\tilde{\Gamma}(\overline{B} \rightarrow D\mu\bar{\nu}_\mu)_{SM}} \quad (\text{including error of } \rho \text{ parameter})$$

$$(\tilde{\Gamma} \rightarrow m_\tau^2 \leq q^2 \leq (m_B - m_D)^2)$$



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$$R = \frac{m_W}{m_H} \tan \beta$$

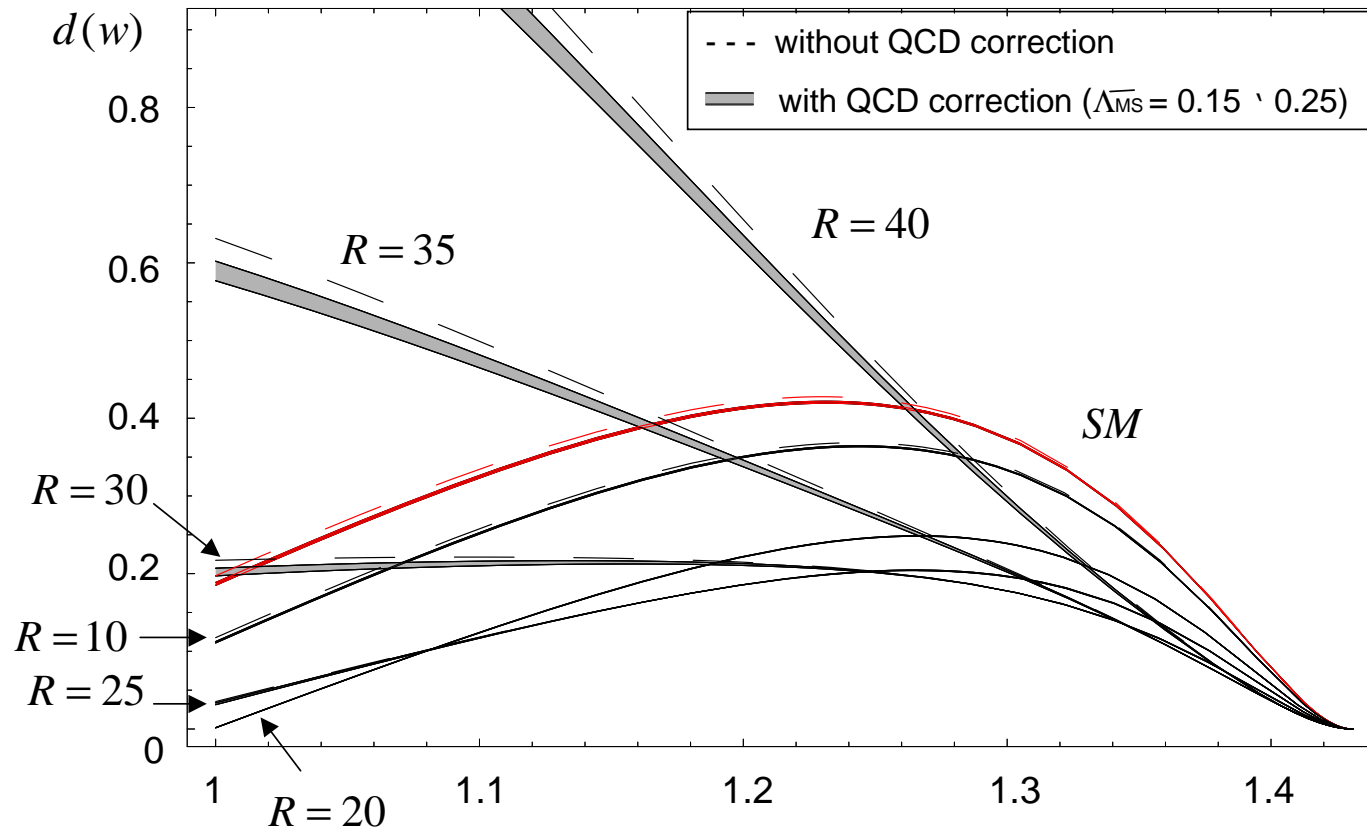
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Distribution

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$$d(w) = (w^2 - 1) \frac{\Gamma(\bar{B} \rightarrow D\tau\bar{\nu}_\tau)/dw}{\Gamma(\bar{B} \rightarrow D\mu\bar{\nu}_\mu)_{SM}/dw}$$

$$R = \frac{m_W}{m_H} \tan \beta$$



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$$w \equiv v \cdot v' = (m_B^2 + m_D^2 - q^2) / 2m_B m_D$$

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5. Summary and discussion

- SUSY effects could be found in the $b \rightarrow s$ transitions. Correlation among several observables would tell us about the SUSY breaking. Hadronic uncertainty should be studied in detail. Expected experimental precision of CP and FB asymmetries?
- The charged Higgs in the MSSM affects BR and the decay distribution of $\bar{B} \rightarrow D\tau\bar{\nu}$. The deviation would be seen only in B factories (not in hadron machines). The hadronic uncertainty is understood well and can be reduced by data from B factories. SUSY radiative correction is significant for large $\tan\beta$ and may enhance the deviation. Can the distribution be measured? How precisely? (Nozaki's talk)

Summary and discussion (cont'd)

Subjects not discussed

- What about $b \rightarrow d$ transitions?

$B \rightarrow X_d \gamma, X_d l^+ l^-, X_d \nu \bar{\nu}, \pi\pi, \rho\pi$, etc.

Δm_{B_d} gives a constraint,

$$|M_{12}(\text{SUSY})| \lesssim |M_{12}(\text{SM})|.$$

The isospin analysis in $\pi\pi$ and $\rho\pi$ is important.

SUSY penguin contributes to the $\Delta I = 1/2$ part.

Further study is necessary.

- Model-independent ϕ_3 measurement

Use tree processes like $B \rightarrow DK$.

ϕ_3 can differ from the value of the CKM fitting in the SM.

(Okada's talk)

- Hadron machines