

原子遷移過程で探る素粒子物理

田中 実 (阪大)

杜の都素粒子ワークショップ —Yasuhiro Okada Festa—, 2025/3/13-14



右から、
岡田さん
田中
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野村さん

“New Initiatives in Muon Lepton Flavor Violation and Neutrino Oscillation
with High Intense Muon and Neutrino Sources”

University of Hawaii in Honolulu, Oct. 2-6, 2000

Introduction

Precision frontier in fundamental physics

Development of precision spectroscopy / optical clock

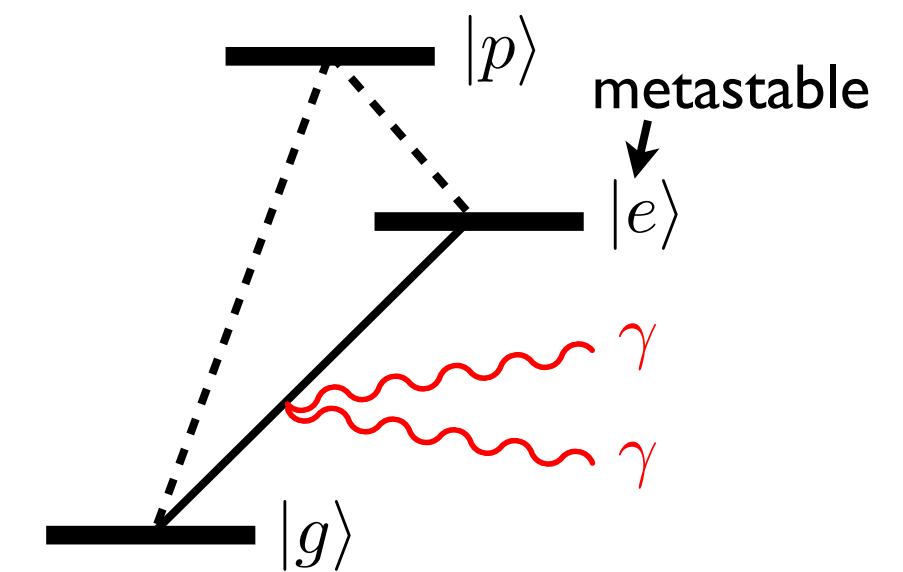
Temporal variation of fundamental constants, α , m_e/m_p

Yb^+ : $\delta\nu/\nu \sim 10^{-18}$, $\delta\nu \sim \text{sub Hz}$ Huntemann et al. (PTB) 2016

Rate enhancement with quantum coherence

Paired Super-Radiance (PSR): $|e\rangle \rightarrow |g\rangle + \gamma + \gamma$

M. Yoshimura, N. Sasao, MT, PRA86, 013812 (2012)



macroscopic amplification by coherence ρ_{eg}

Para-hydrogen PSR experiment @Okayama U

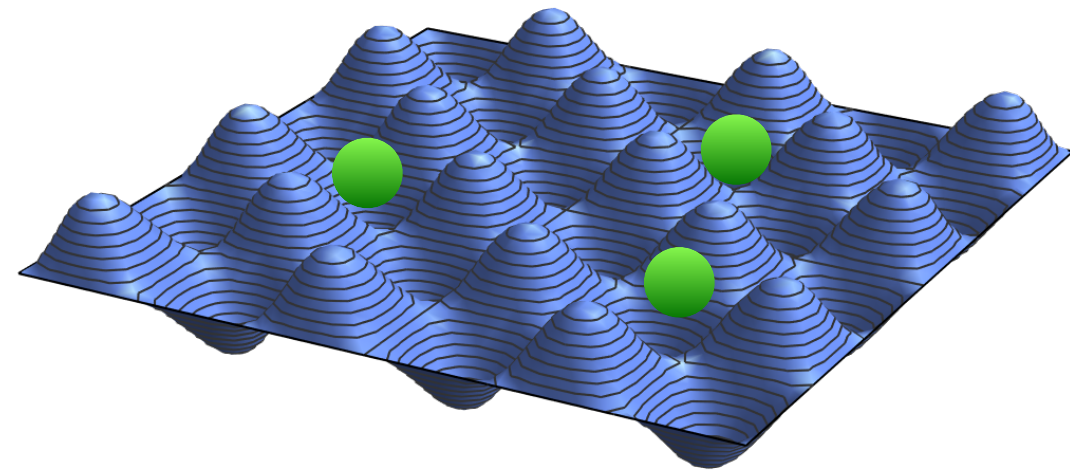
Y. Miyamoto et al. PTEPI 13C01 (2014), PTEP081C01 (2015)

10^{18} enhancement achieved

今, 主にやっていること

Isotope shift with high precision spectroscopy

search for new neutron-electron interaction

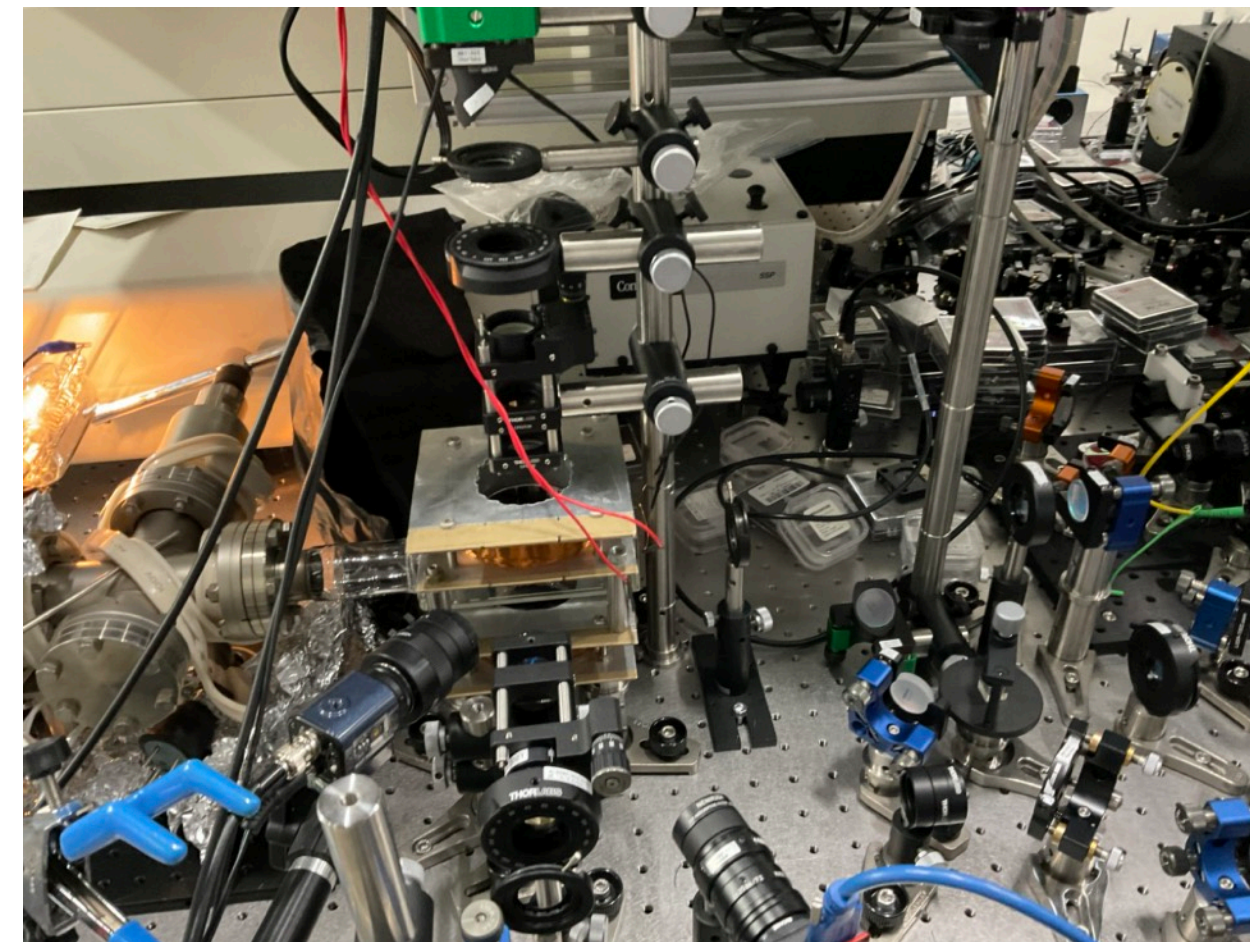


Yb atoms in an optical lattice

Takahashi group, Kyoto U

Dark matter search with macro-coherent amplification

light halo dark matter,
e.g. axion, dark photon



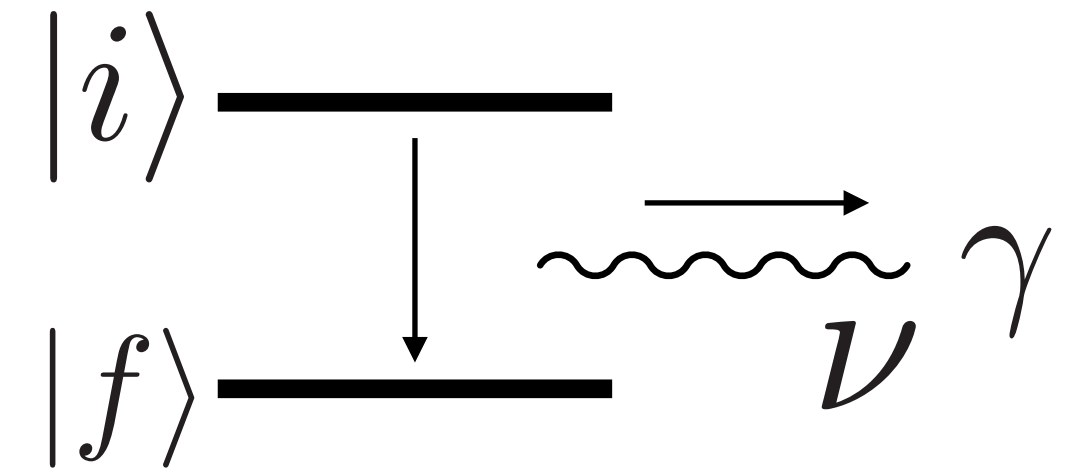
Cesium vapor
@Okayama U

Isotope Shift

Isotope shift (IS) and King linearity

Level-splitting difference between isotopes

$$IS = \nu_{A'A} := \nu_{A'} - \nu_A$$



IS of two transitions: $t = 1, 2$

$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A}$$

$$\mu_{A'A} := 1/m_{A'} - 1/m_A$$

$$\langle r^2 \rangle_{A'A} := \langle r^2 \rangle_{A'} - \langle r^2 \rangle_A$$

mass shift (MS) field shift (FS)

Modified IS: $\tilde{\nu}_{A'A}^{(t)} := \nu_{A'A}^{(t)} / \mu_{A'A} = \boxed{K_t} + \boxed{F_t \langle r^2 \rangle_{A'A} / \mu_{A'A}}$

electronic factors

nuclear factor

King linearity: eliminating the nuclear factor King, 1963

$$\tilde{\nu}_{A'A}^{(2)} = K_{21} + F_{21} \tilde{\nu}_{A'A}^{(1)}$$

$$K_{21} := K_2 - F_{21} K_1, \quad F_{21} := F_2 / F_1$$

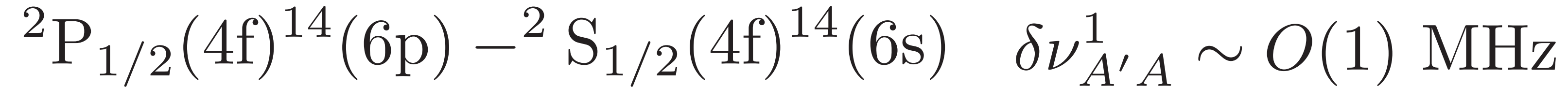
→ $(\tilde{\nu}_{A'A}^{(1)}, \tilde{\nu}_{A'A}^{(2)})$ on a straight line, King plot

Ex. Yb⁺

K. Mikami, MT, Y. Yamamoto EPJC77:896 (2017)

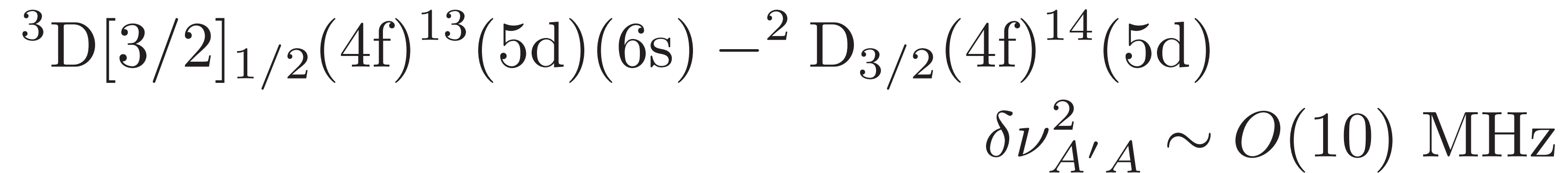
Transition 1: 369 nm

Martensson-Pendrill et al. PRA49, 3351 (1994)



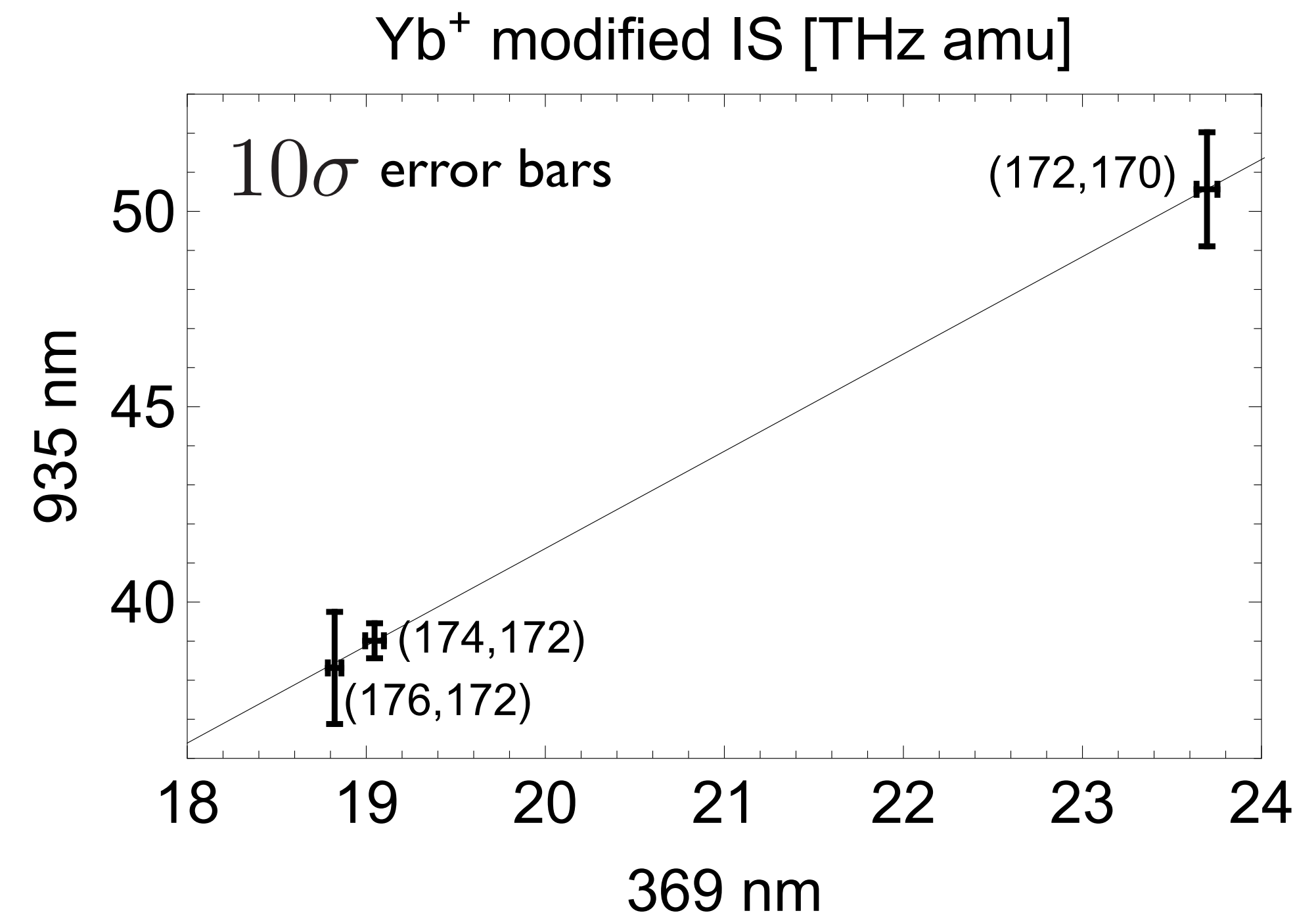
Transition 2: 935 nm

Sugiyama et al. CPEM2000



Isotope pairs

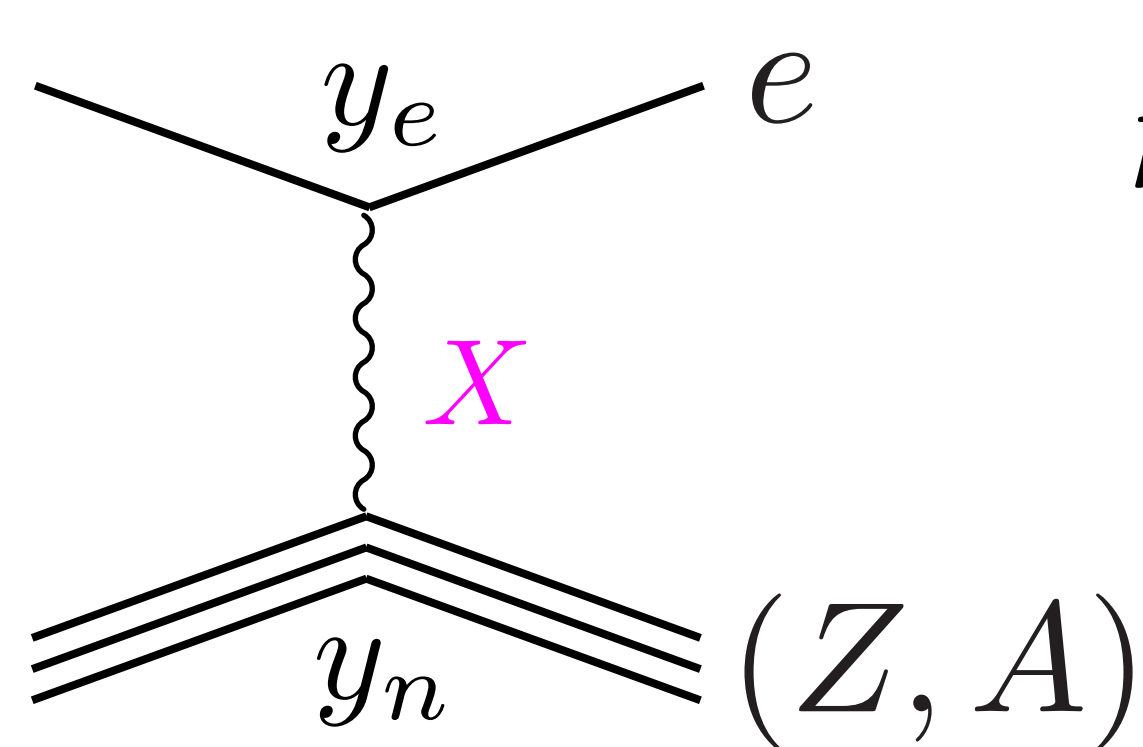
(172, 170), (174, 172), (176, 172)



Nonlinearity

IS by new **neutron-electron interaction**

Delaunay et al. arXiv:1601.05087v2



$$\nu_{A'A}^{(t)} = \underbrace{K_t \mu_{A'A}}_{\text{MS}} + \underbrace{F_t \langle r^2 \rangle_{A'A}}_{\text{FS}} + \underbrace{X_t (A' - A)}_{\text{particle shift (PS)}}$$

Nonlinearity due to **subleading FS**

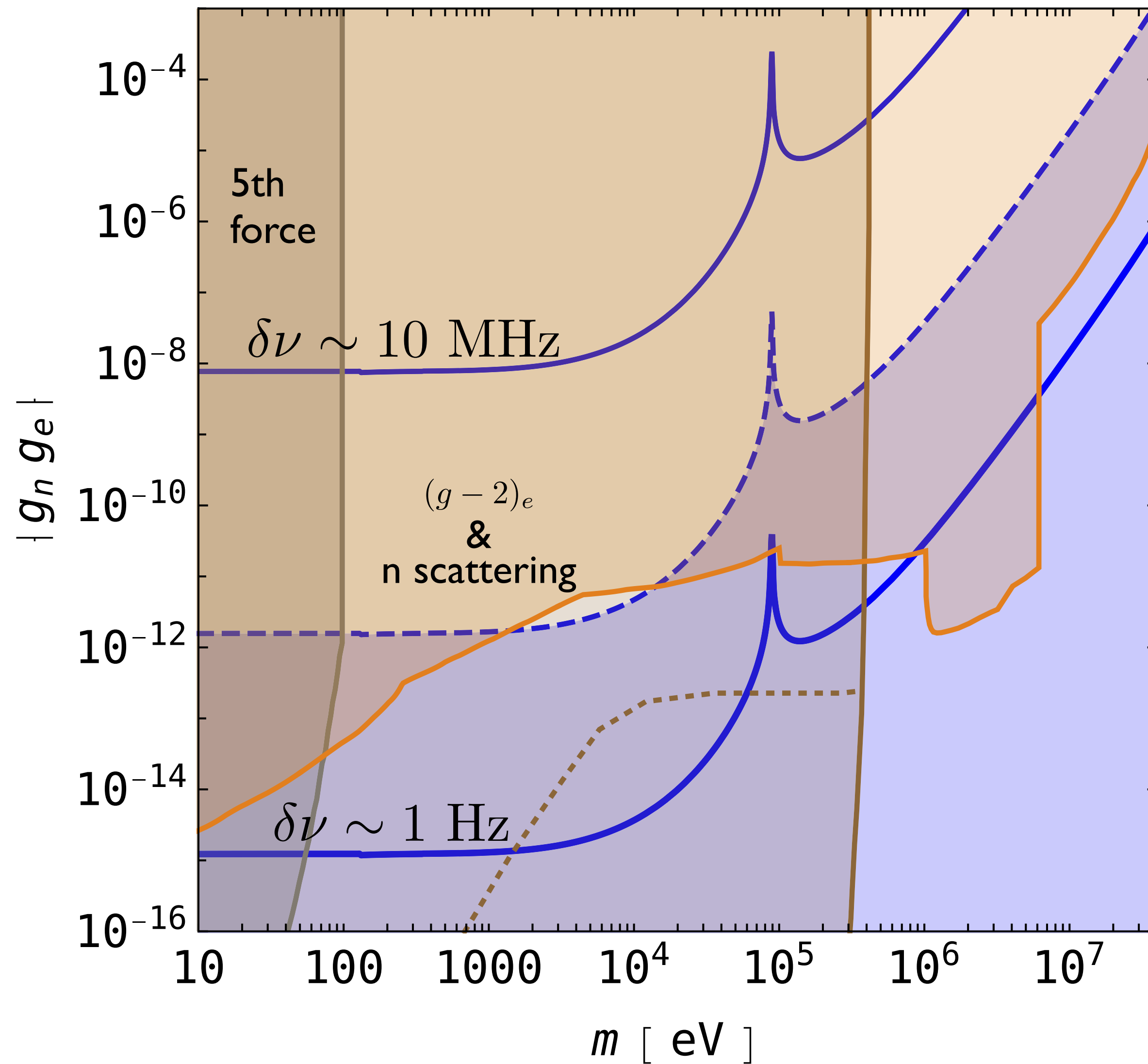
$$\text{FS} = F_t \langle r^2 \rangle_{A'A} + F'_t [\langle r^2 \rangle_{A'A}]^2 + G_t \langle r^4 \rangle_{A'A} + \dots$$

quadratic FS (FS₂₂) higher moment (FS₄)

$$[\langle r^2 \rangle_{A'A}]^2 := (\langle r^2 \rangle_{A'A_0})^2 - (\langle r^2 \rangle_{AA_0})^2$$

Ex. Yb⁺

MT, Y. Yamamoto PTEP I03B02 (2020)



Transition 1: 369 nm Martensson-Pendrill et al. PRA49, 3351 (1994)
 $^2P_{1/2}(4f)^{14}(6p) - ^2S_{1/2}(4f)^{14}(6s)$ $\delta\nu_{A'A}^1 \sim O(1)$ MHz

Transition 2: 935 nm Sugiyama et al. CPEM2000
 $^3D[3/2]_{1/2}(4f)^{13}(5d)(6s) - ^2D_{3/2}(4f)^{14}(5d)$
 $\delta\nu_{A'A}^2 \sim O(10)$ MHz

— **Yb⁺ bounds**
 - - - **$\langle r^4 \rangle$ FS nonlinearity (SM BG)**

FSNL dominance:

$$\delta\nu \lesssim 1 \text{ kHz}$$

What about SM nonlinearity?

Precise calculation difficult

Generalized linearity

K. Mikami, MT, Y. Yamamoto EPJC77:896 (2017)

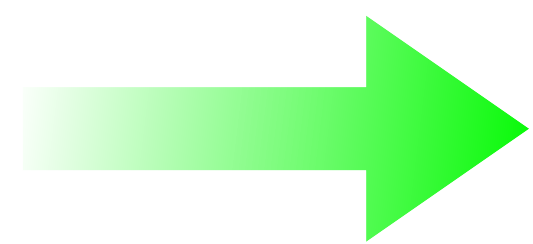
$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A} + F'_t [\langle r^2 \rangle_{A'A}]^2 + X_t (A' - A)$$

3 transitions: $t=1, 2, 3$

FS₂₂

PS

$$\begin{pmatrix} \nu_{A'A}^{(1)} - X_1(A' - A) \\ \nu_{A'A}^{(2)} - X_2(A' - A) \\ \nu_{A'A}^{(3)} - X_3(A' - A) \end{pmatrix} = \begin{pmatrix} K_1 & F_1 & F'_1 \\ K_2 & F_2 & F'_2 \\ K_3 & F_3 & F'_3 \end{pmatrix} \begin{pmatrix} \mu_{A'A} \\ \langle r^2 \rangle_{A'A} \\ [\langle r^2 \rangle_{A'A}]^2 \end{pmatrix} =: M \begin{pmatrix} \mu_{A'A} \\ \langle r^2 \rangle_{A'A} \\ [\langle r^2 \rangle_{A'A}]^2 \end{pmatrix}$$



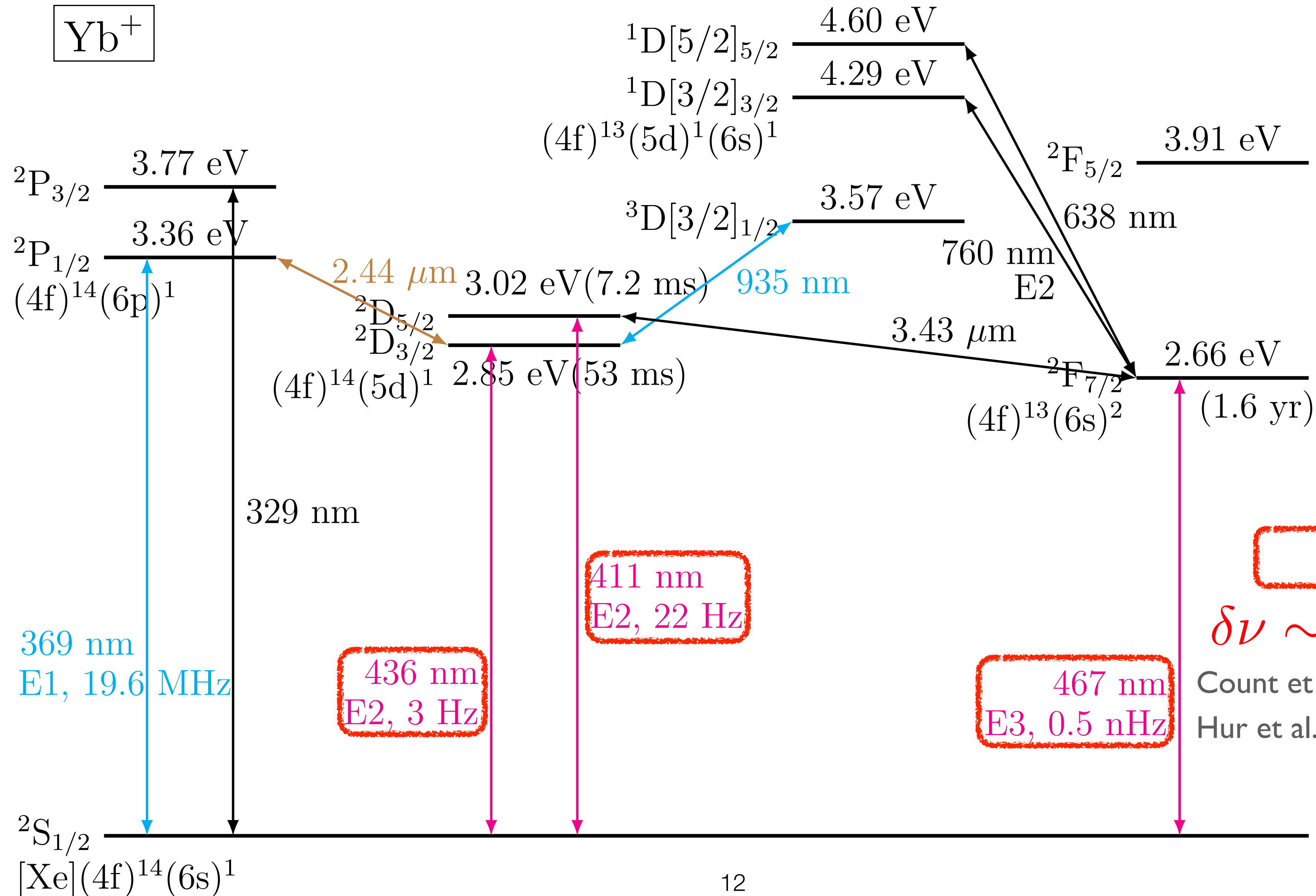
$$\begin{aligned} & (M^{-1})_{11} \nu_{A'A}^{(1)} + (M^{-1})_{12} \nu_{A'A}^{(2)} + (M^{-1})_{13} \nu_{A'A}^{(3)} \\ & - \{ (M^{-1})_{11} X_1 + (M^{-1})_{12} X_2 + (M^{-1})_{13} X_3 \} (A' - A) = \mu_{A'A} \end{aligned}$$

$$(\nu_{A'A}^{(1)}, \nu_{A'A}^{(2)}, \nu_{A'A}^{(3)}) / \mu_{A'A}$$

on a plane if $X_t = 0$

n transitions and $n+1$ IS pairs \rightarrow NP search with $n-2$ NL's removed

Data of Yb transitions



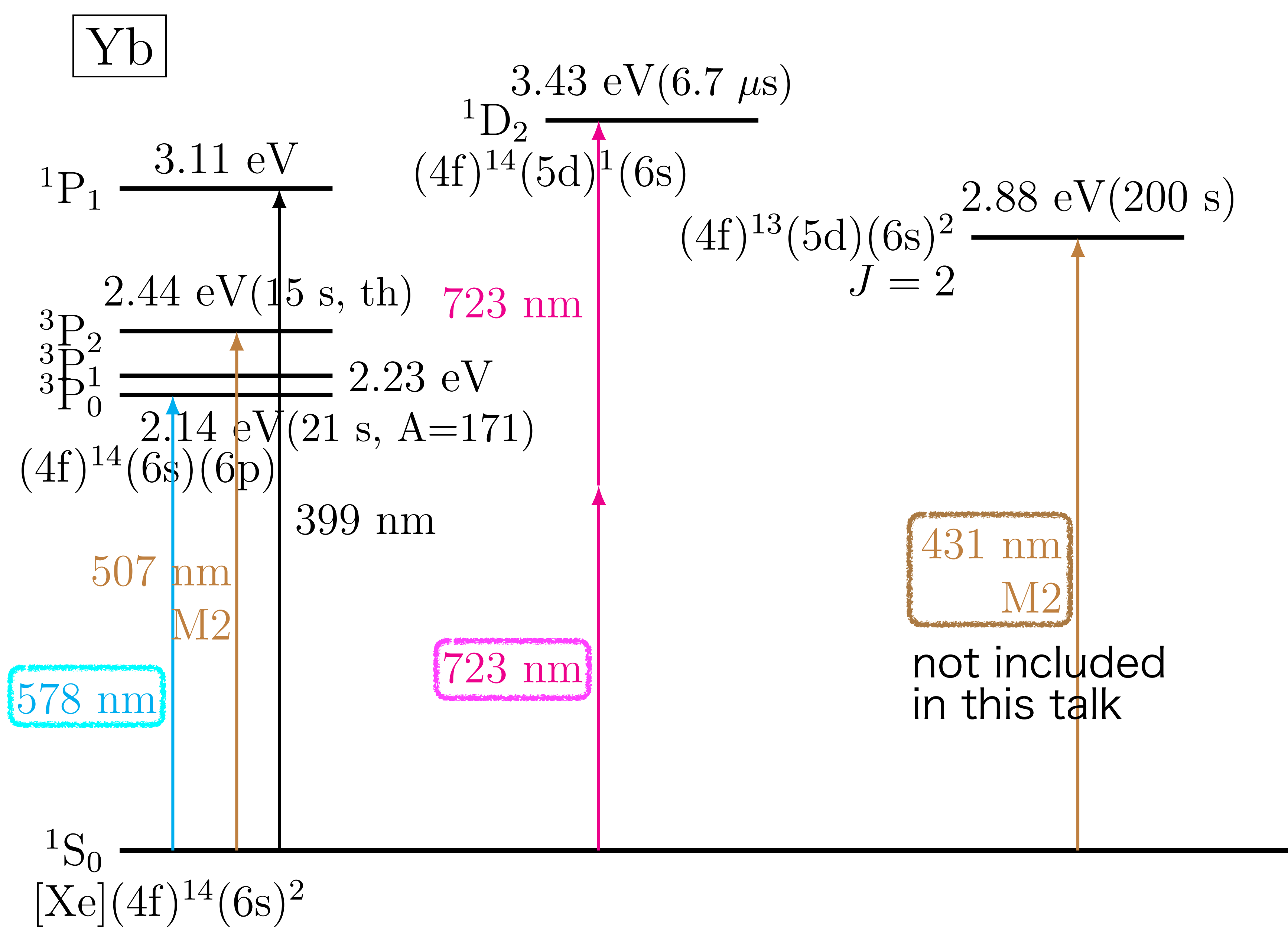
LS coupling
 $2s+1 L_J$

MIT

$\delta\nu \sim 300 - 500 \text{ Hz}$

Count et al. PRL 125, 123002 (2020)

Hur et al. PRL 128, 163201 (2022)



LS coupling

$$2s+1 L_J$$

Kyoto

$\delta\nu \sim$ a few Hz

K. Ono, MT et al. PRX 12, 021033 (2022)

Mainz

Figueroa et al. PRL 128, 073001 (2022)

$\delta\nu \sim O(100)$ Hz

Kyoto

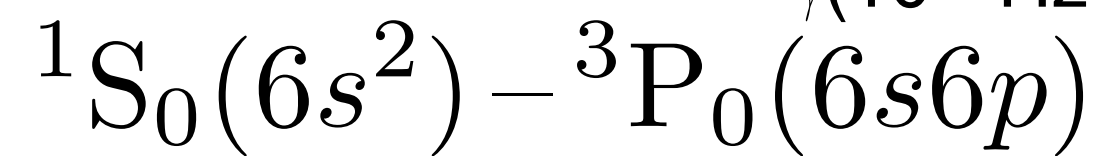
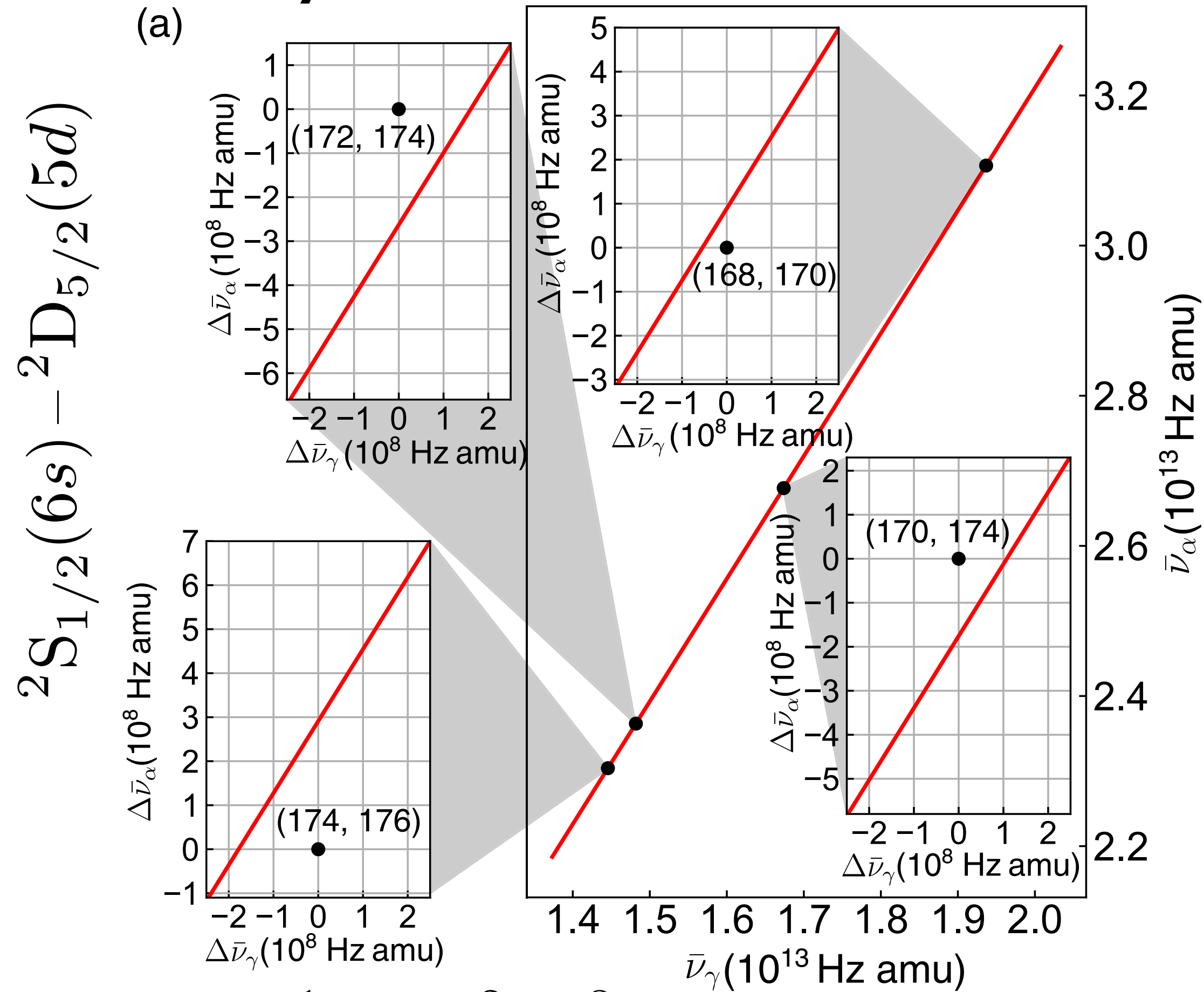
T. Ishiyama et al. PRL 130, 153402 (2023)

NMIJ, AIST

A. Kawasaki et al. PRA 109, 062806 (2024)

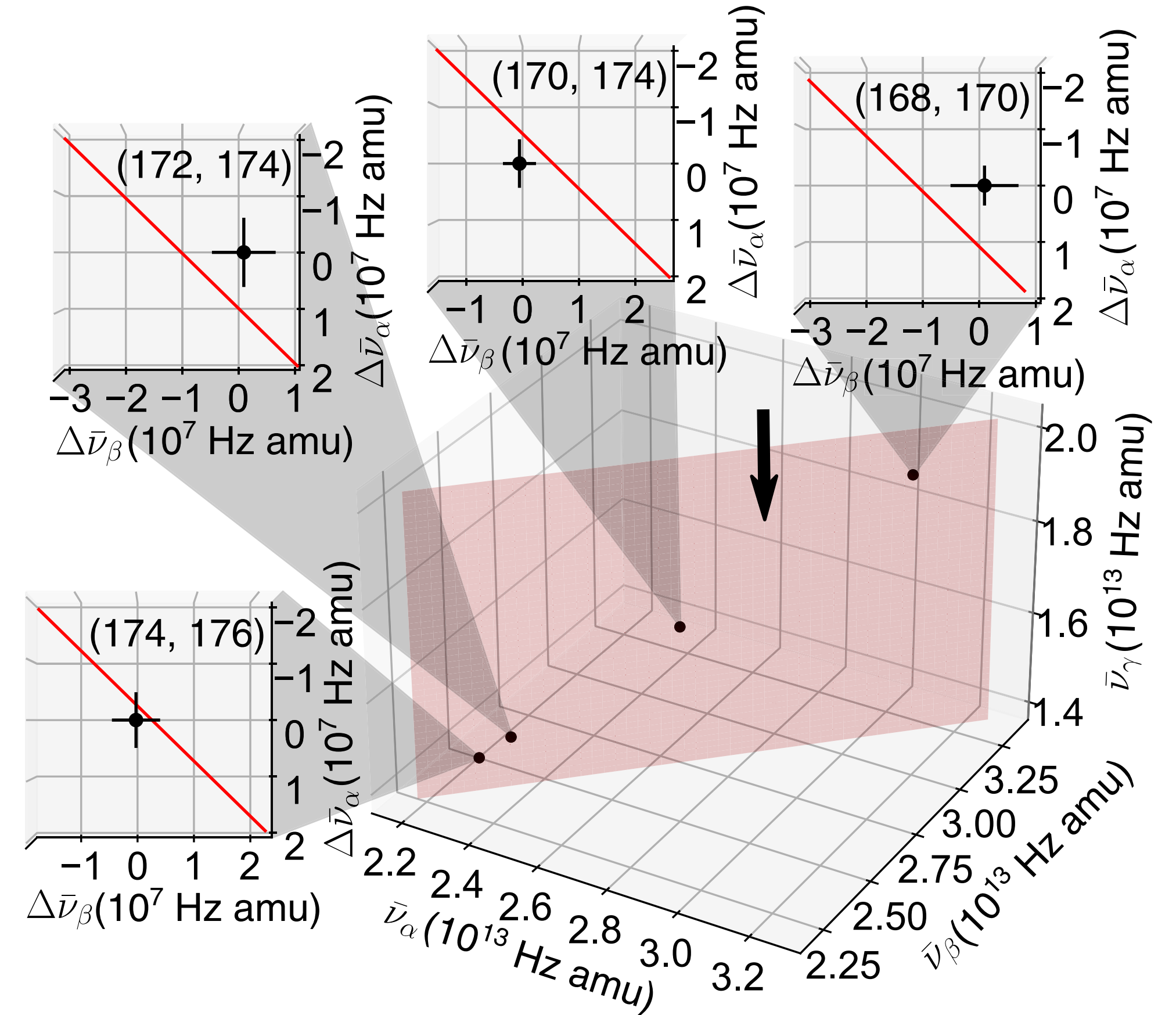
431 nm
M2
not included
in this talk

2D analysis



$\chi^2/\text{dof} = 1.1 \times 10^4 / 3$

3D analysis



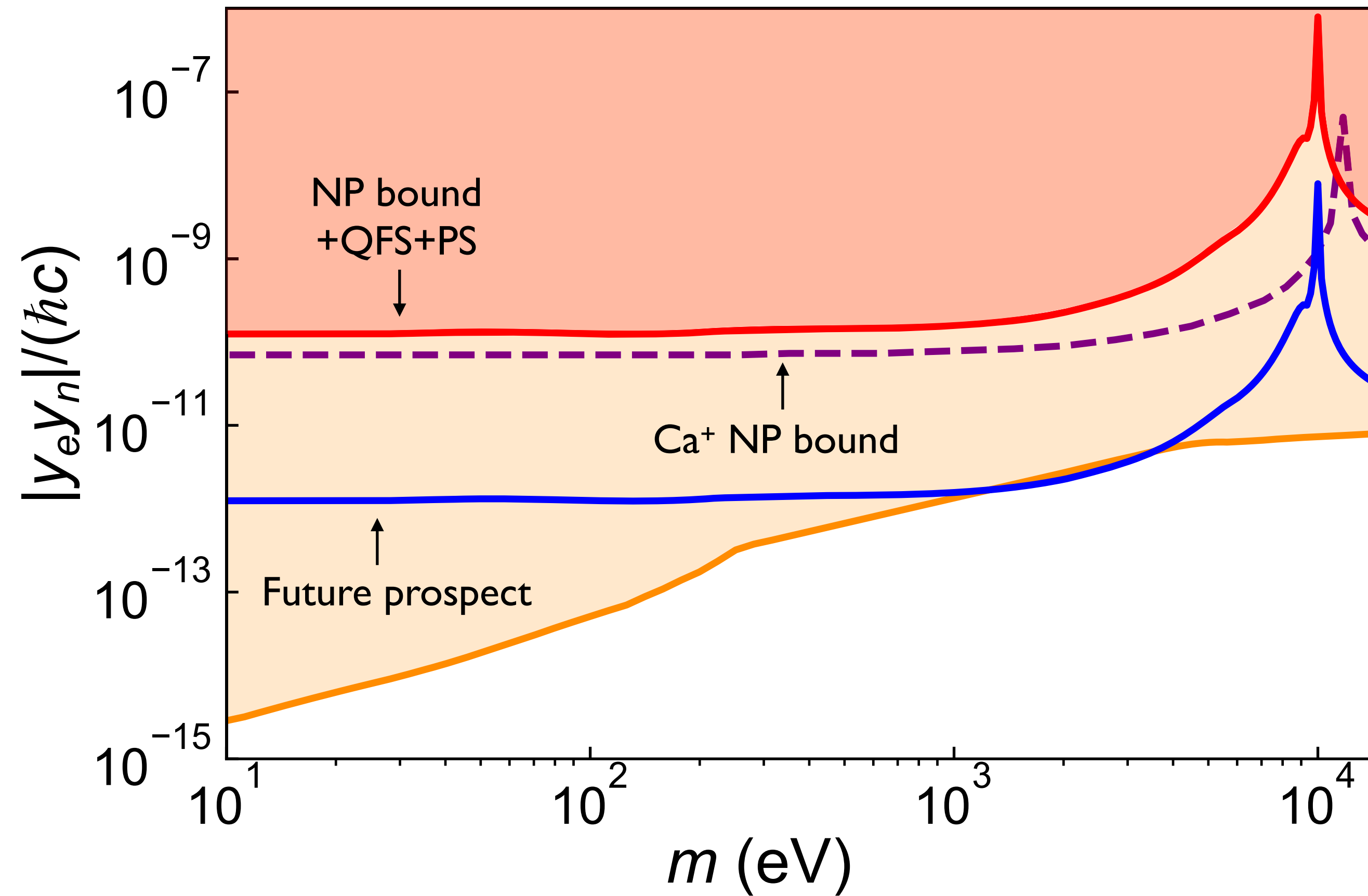
$\chi^2/\text{dof} = 15/3$



two or more NL sources

+QFS+PS \longrightarrow New physics bound

K. Ono, MT et al. PRX 12, 021033 (2022)



Simultaneous analysis of plural linearities

Yb: 4 independent IS pairs

3D linearity is nontrivial, while 5 transitions available.

→ Combined fit of 3 independent 3D relations

$$\nu_3 = k_3\mu + f_{31}\nu_1 + f_{32}\nu_2, \nu_4 = k_4\mu + f_{41}\nu_1 + f_{42}\nu_2, \nu_5 = k_5\mu + f_{51}\nu_1 + f_{52}\nu_2$$

(A'A omitted)

$$\delta\chi^2/\text{dof} = 35.1/3$$

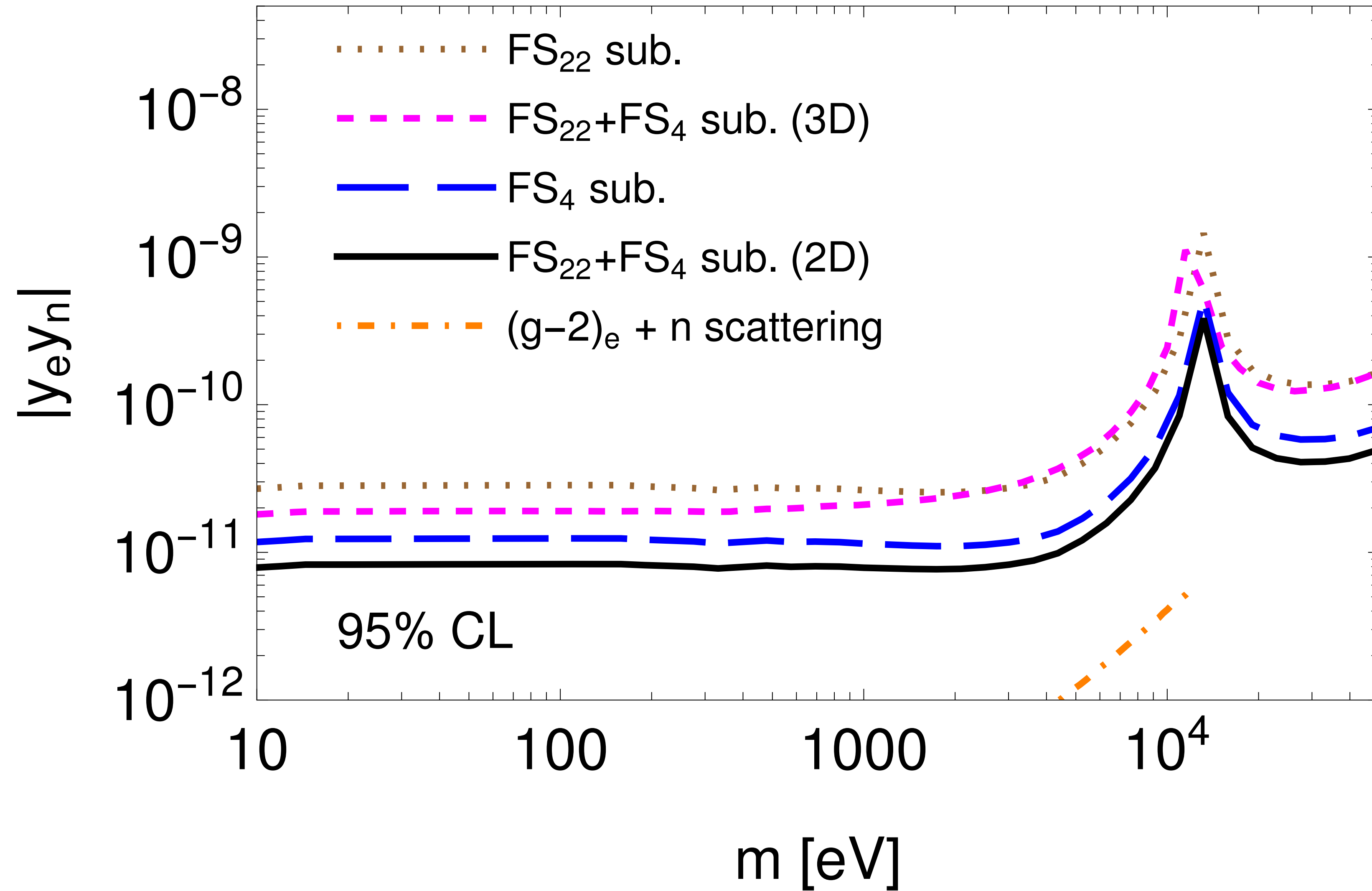
Assigning this nonlinearity to PS,

→ $\delta\chi^2/\text{dof} = 9.98/1$

PS alone cannot explain the observed 3D nonlinearity.

Five transitions combined

MT, Y. Yamamoto @JPS meeting, Mar. 2024



Dark Matter Search

Light dark matter candidates

Axion (and axion-like particles)

Proposed solution of the strong CP problem

$$\mathcal{L}_{\text{int}} \supset g_{a\gamma\gamma} a F^{\mu\nu} \tilde{F}_{\mu\nu} \propto a \mathbf{E} \cdot \mathbf{B}$$

2 parameters : $m_a, g_{a\gamma\gamma}$

Dark photon

Minimal extension of the SM gauge sector $U_X(1)$

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \boxed{\frac{\chi}{2} F_{\mu\nu} X^{\mu\nu}} + \frac{1}{2} m_X^2 X_\mu X^\mu - j_{\text{em}}^\mu A_\mu$$

2 parameters : m_X, χ

kinetic mixing

Standard model + inflation + dark photon

Graham, Mardon, Rajendran, PRD93, 103520 (2016)

➡ Dark photon dark matter

$$\Omega_X = \Omega_{\text{CDM}} \left(\frac{m_X}{6 \times 10^{-6} \text{ eV}} \right)^{1/2} \left(\frac{H_I}{10^{14} \text{ GeV}} \right)^2$$

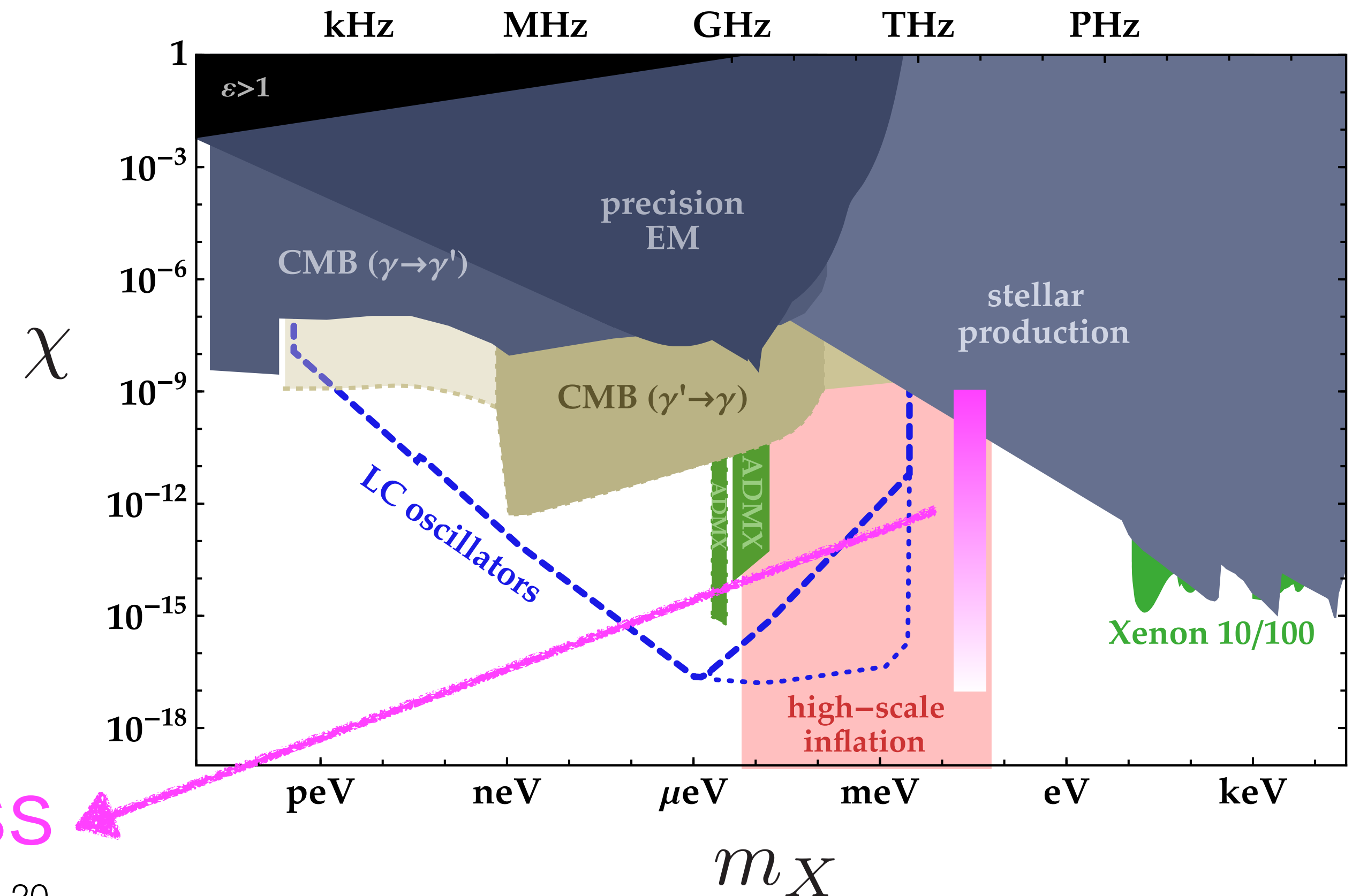
H_I : Hubble scale of inflation

$$H_I \lesssim 10^{14} \text{ GeV}$$

➡ $m \gtrsim 10^{-5} \text{ eV}$

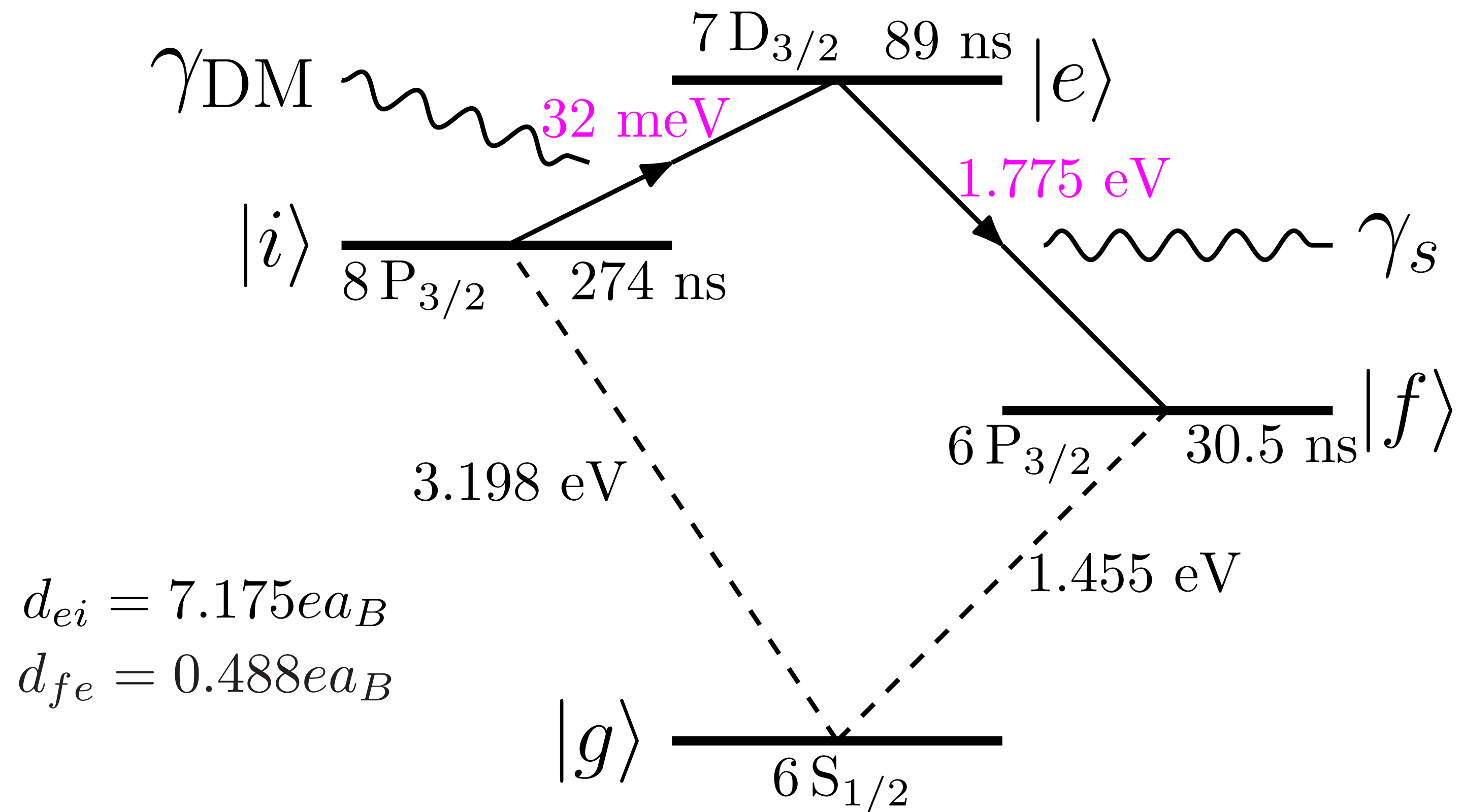
coherently oscillating

$$X^\mu \propto e^{im_X t}$$



Search with atomic process

Cs pilot experiment @Okayama U



$$d_{ei} = 7.175ea_B$$

$$d_{fe} = 0.488ea_B$$

target spec.

$$n = 1 \times 10^{12} \text{ cm}^{-3}$$

$$V = 0.1 \times 0.1 \times 1 \text{ cm}^3$$

$$\Gamma_{\text{eff}} \simeq 2.2 \times 10^{-5} \text{ eV}$$

(\gg natural width)

$$\Gamma = 1.6 \times 10^3 \left(\frac{\chi}{10^{-9}} \right)^2 \left(\frac{n}{10^{12} \text{ cm}^3} \right)^2 \left(\frac{\rho_{fi}}{0.25} \right)^2 \text{ Hz}$$

cf. single atom rate: $\Gamma_0 = 3.7 \times 10^{-9} \left(\frac{\chi}{10^{-9}} \right)^2 \text{ Hz}$

Coherence generation

Liouville - von Neumann equation with relaxation

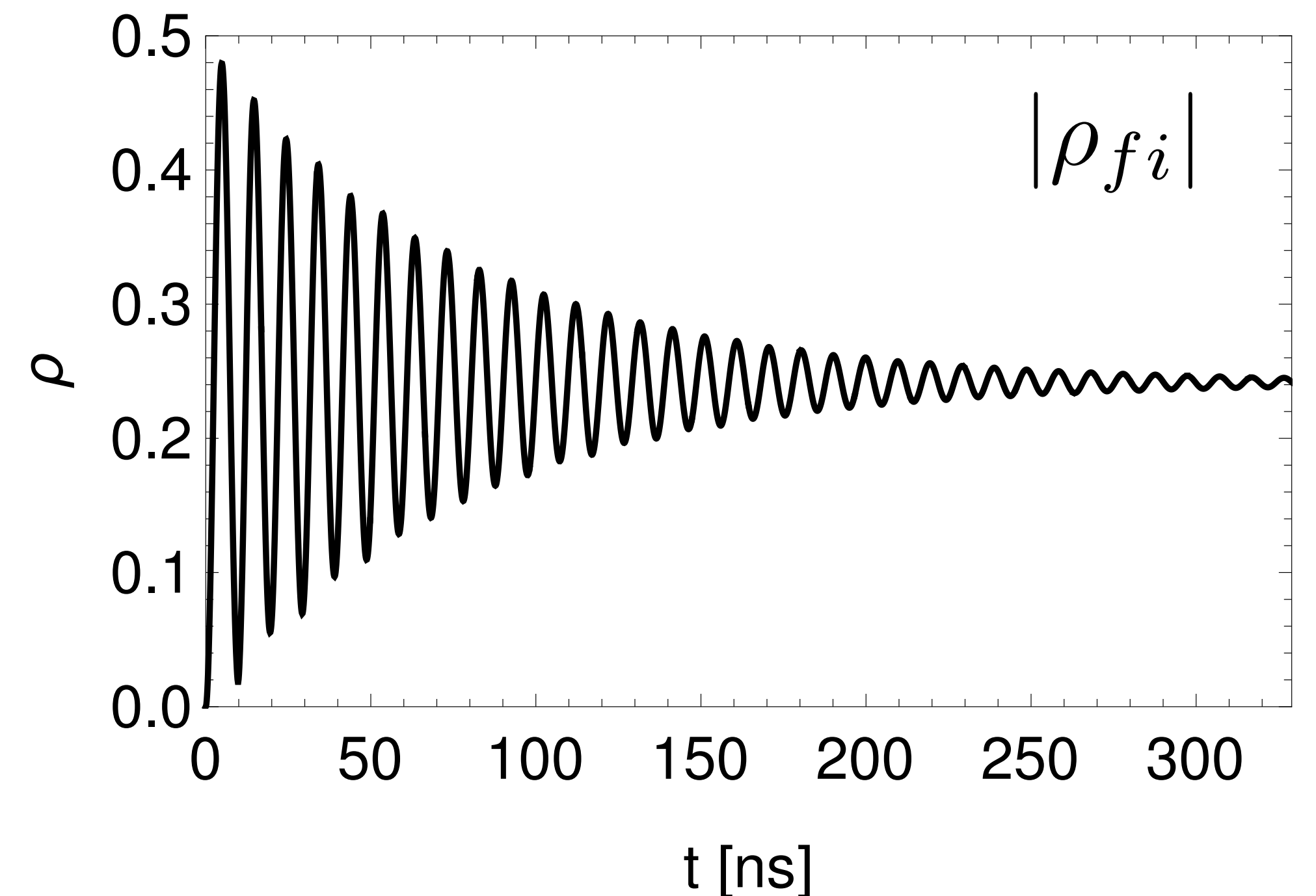
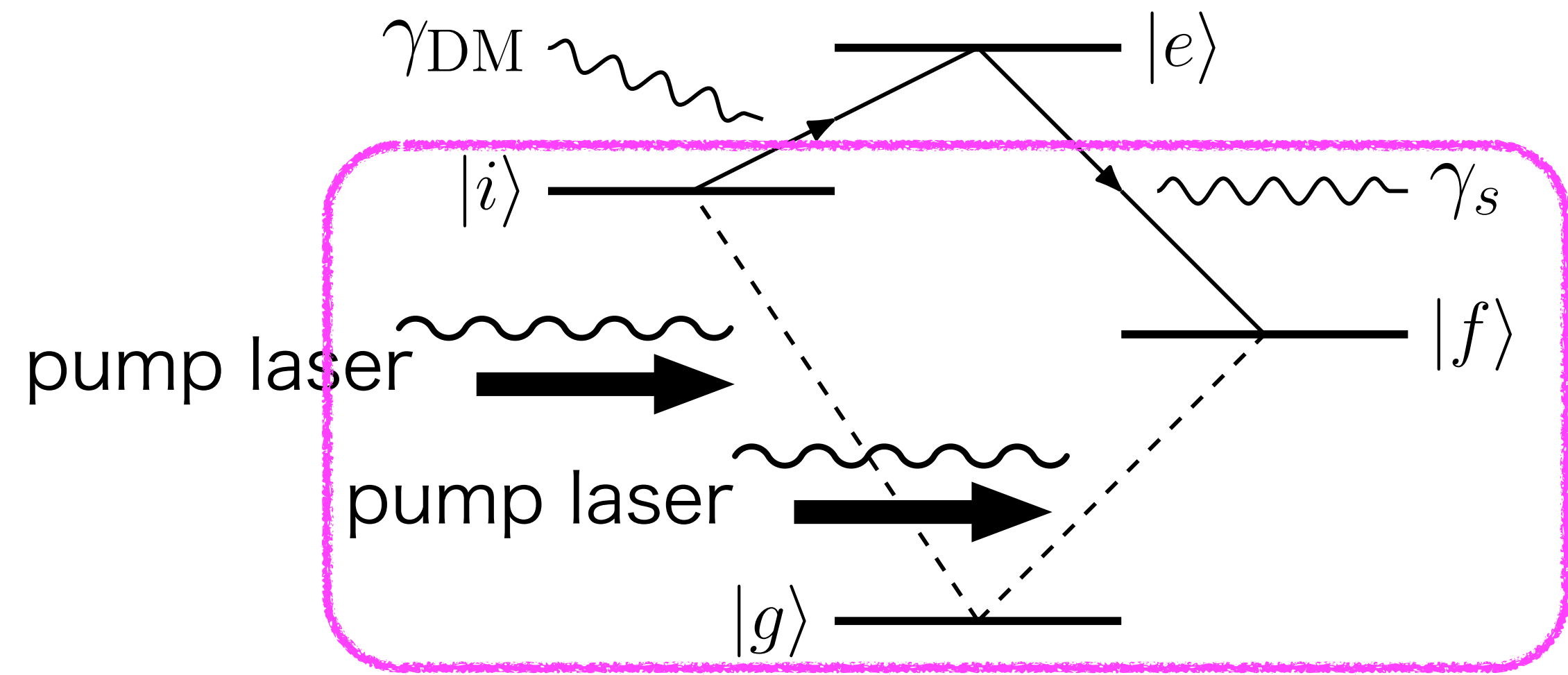
$$\partial_t \rho_{ij}(t) = -i[H(t), \rho(t)]_{ij} - \sum_{m,n} \Gamma_{ij,mn} \rho_{mn}$$

Cs parameters with radiation damping only

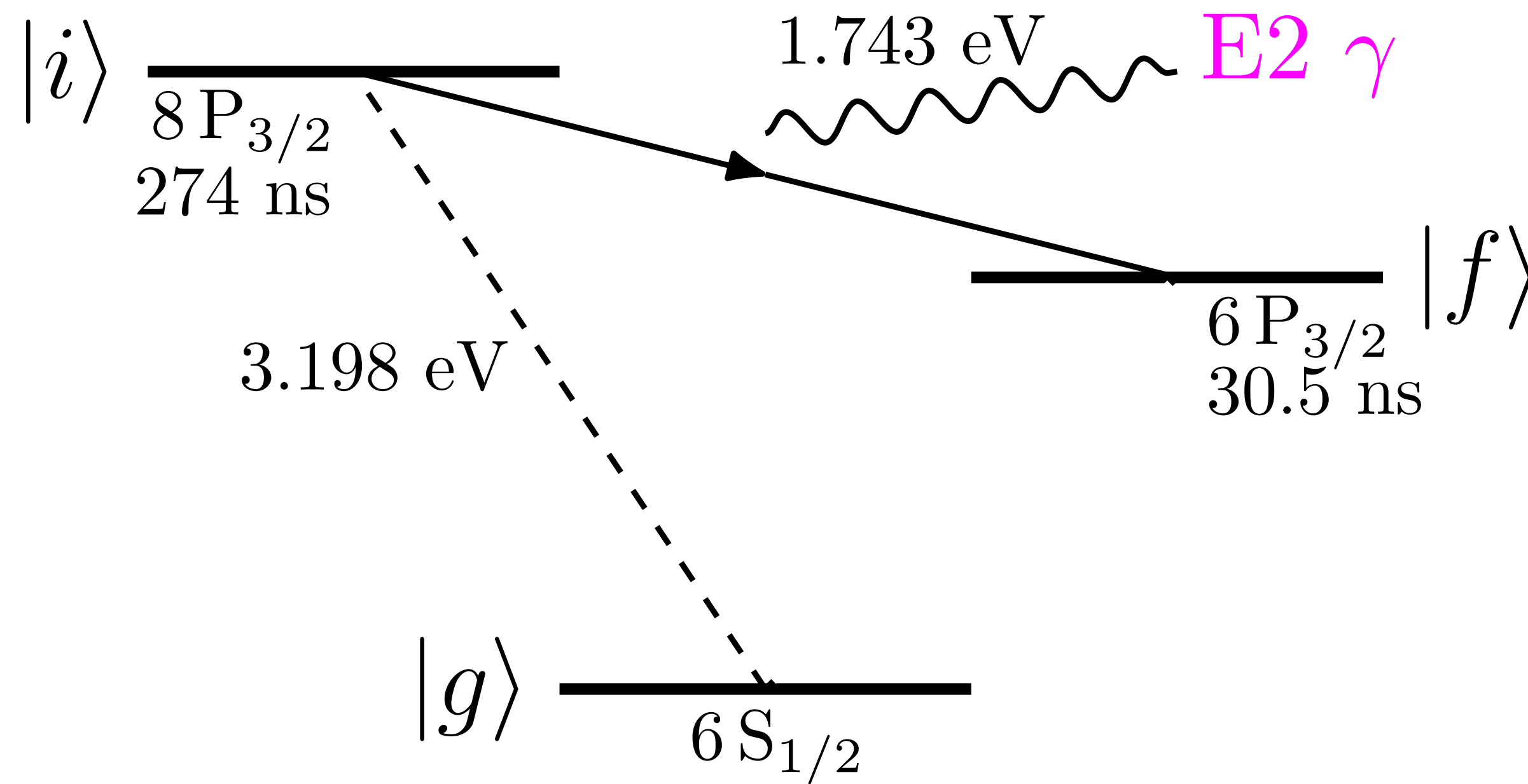
CW laser power: 1W(g-i), 1mW(g-f)

Laser cross section: 1mm²

$|\rho_{fi}| \simeq 1/4$ possible



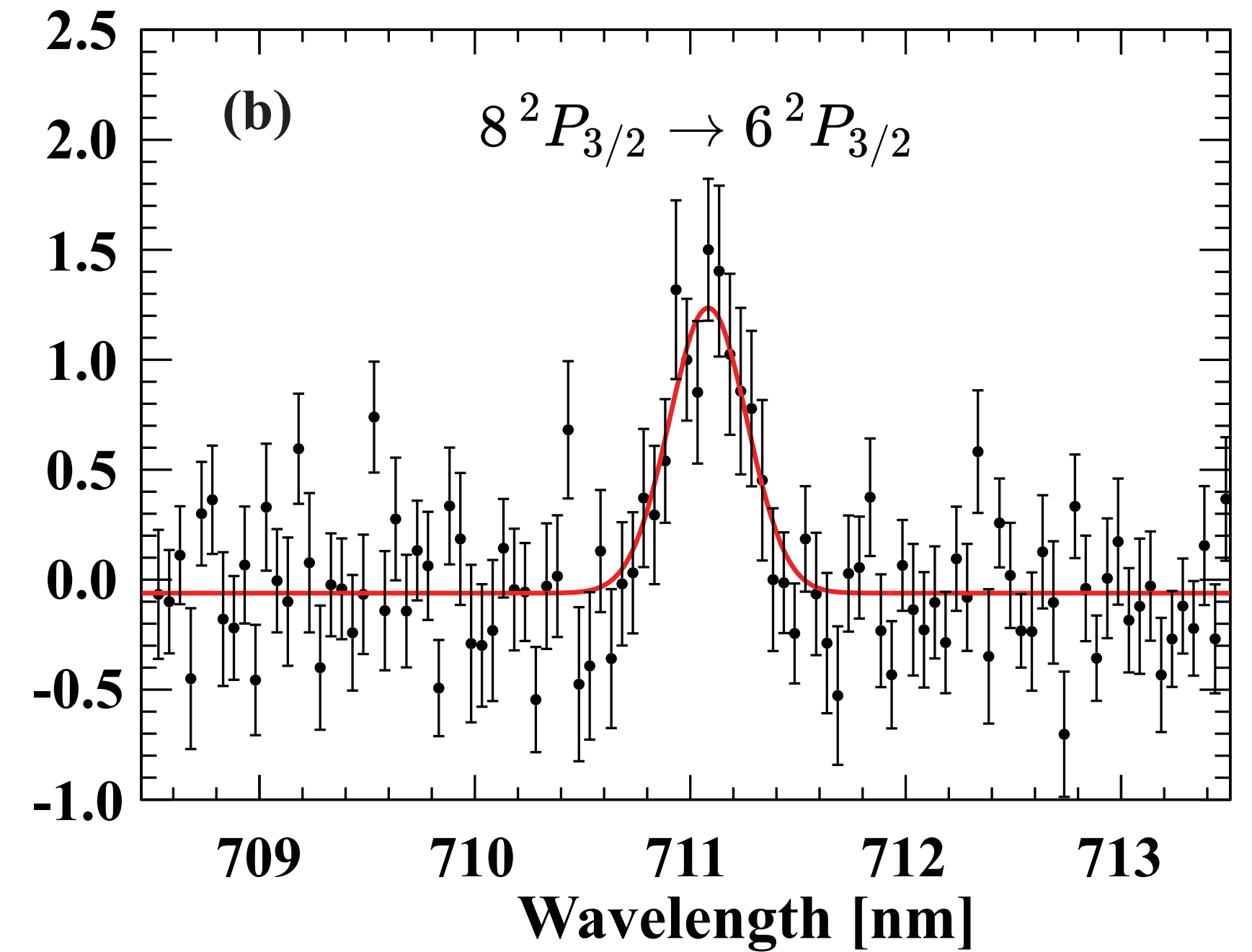
Coherence measurement



E2 transition

$$\Gamma(i \rightarrow f) \sim \text{a few Hz}$$

J. Wang, MT et al. PRA 110, 012804 (2024)



Observed w/o coherence

Rate measurement with coherence $\longrightarrow \rho_{fi}$

applying the second pump laser, ongoing

Summary

Isotope shift

- **Isotope shift and King linearity** $\tilde{\nu}_{A'A}^{(2)} = K_{21} + F_{21}\tilde{\nu}_{A'A}^{(1)}$

IS=MS+FS, linear relation of mIS of two transitions

- **Nonlinearities**: New physics and/or SM higher order

- **Generalized linearity**

SM nonlinearity removed, improved sensitivity to new physics

Two or more SM higher order contributions revealed

One SM higher order + PS, excluded by our combined analysis

- **New precise Yb IS data**

Yb⁺ ion O(10) Hz, MPI-PTB Door et al., PRL 134 063002 (2025)

improvement of MIT data and Yb masses

Yb atom <O(10) Hz, Kyoto

$^1S_0(6s)^2 - (4f)^{13}(5d)(6p)^2$ ($J = 2$) 431 nm, $^1S_0(6s)^2 - ^3P_2(6s)(6p)$ 507 nm

Dark matter search

- Rate amplification by coherence in ensemble of atoms
macrocoherence and momentum conservation
significance of ρ_{fi}
- Halo dark photon/axion search experiment
Cs pilot experiment at Okayama U
coherence generation and ρ_{fi} measurement ongoing
- Theoretical issues
more realistic simulation of signal and background