



# Isotope shift as a probe of new physics

Minoru Tanaka (Osaka U.)

In collaboration with

Y. Yamamoto (NCTS), K. Ono(Kyoto), T. Higomoto(Kyoto), Y. Saito(Kyoto),  
T. Ishiyama(Kyoto), Y. Takasu(Kyoto), T. Takano (Kyoto), Y. Takahashi(Kyoto)

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# Frontiers in particle physics

Energy frontier: LHC, ILC, FCC, ...

Intensity frontier: B factory, K, muon, ...

Cosmic frontier: CMB, GW, ...

Precision / low energy frontier

$0\nu\beta\beta$ , DM, EDM, ...

Temporal variation of fundamental constants

$\alpha$ ,  $m_e/m_p$  using atomic clock  $\text{Yb}^+ : \delta\nu/\nu \sim 10^{-18}$ ,  $\delta\nu \sim \text{sub Hz}$

Huntemann et al. (PTB) 2016

Isotope shift new neutron-electron interaction

# Light new particle search

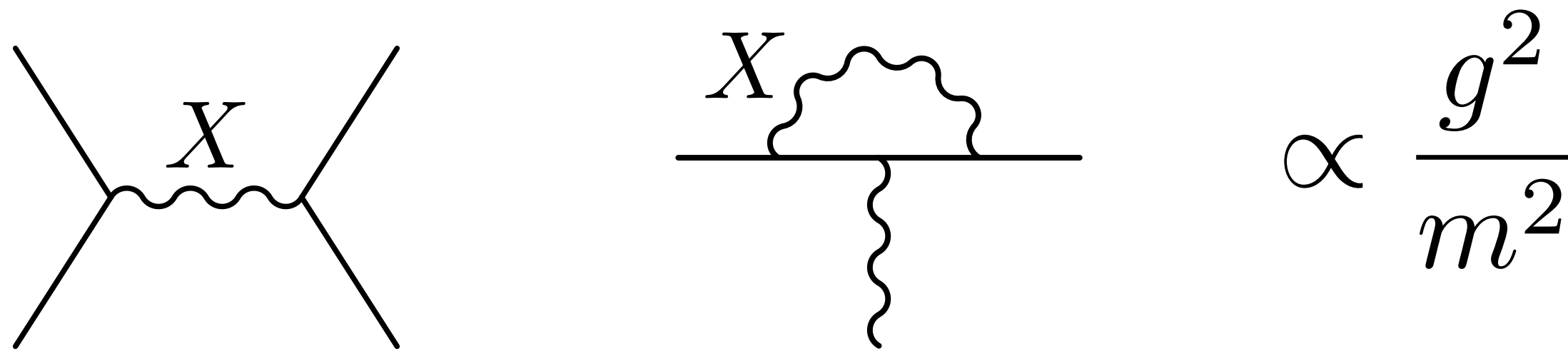
## Direct search

Visible decay, e.g.  $X \rightarrow e^+ e^-$  : direct search

Invisible decay, e.g.  $X \rightarrow \nu \bar{\nu}$  : missing E/p

Stable: missing E/p, dark matter?

## Indirect search



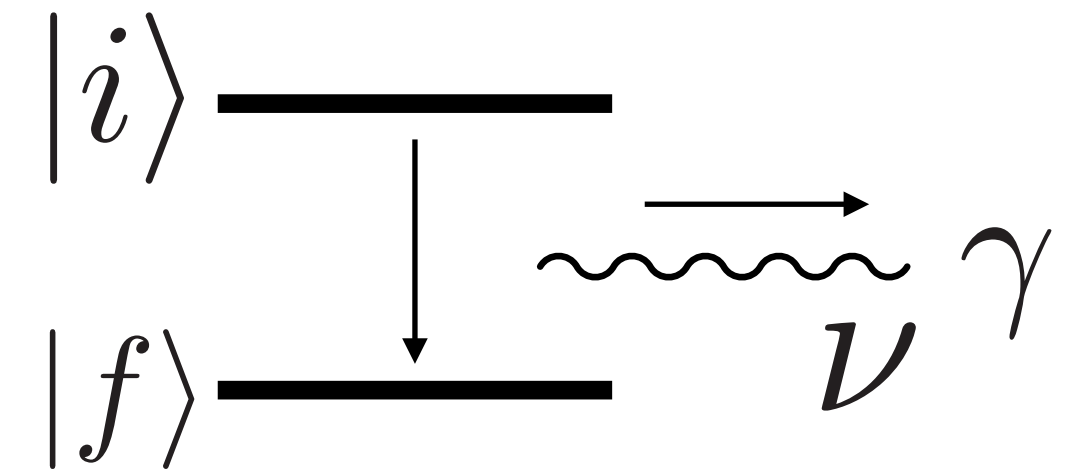
cf. weak interaction  $\sim \frac{g_Z^2}{m_Z^2} \sim \frac{0.5}{(100 \text{ GeV})^2} = \frac{0.5 \times 10^{-10}}{(1 \text{ MeV})^2}$

# Isotope shift (IS)

Level-splitting difference between isotopes

$$h\nu_A = E_A^i - E_A^f, \quad h\nu_{A'} = E_{A'}^i - E_{A'}^f$$

$$\text{IS} = \nu_{A'A} := \nu_{A'} - \nu_A$$



No IS for infinitely heavy and point-like nuclei  $\longrightarrow$   $\text{IS} = \text{MS} + \text{FS}$

Mass shift: finite mass of nuclei (reduced mass)

$$\text{MS} \propto 1/m_{A'} - 1/m_A \quad (\text{dominant for } Z < 20)$$

Field shift: finite size of nuclei

$$\text{FS} \propto \langle r^2 \rangle_{A'} - \langle r^2 \rangle_A \quad (\text{dominant for } Z > 40)$$

Theoretical calculation of IS: Not easy  $\text{IS} \sim O(\text{GHz}) \sim O(10 \mu\text{eV})$



# King linearity

King, 1963

IS of two transitions:  $t = 1, 2$

$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A}$$

mass shift (MS)    field shift (FS)

$$\mu_{A'A} := 1/m_{A'} - 1/m_A$$

$$\langle r^2 \rangle_{A'A} := \langle r^2 \rangle_{A'} - \langle r^2 \rangle_A$$

Modified IS:  $\tilde{\nu}_{A'A}^{(t)} := \nu_{A'A}^{(t)} / \mu_{A'A} = \boxed{K_t} + \boxed{F_t \langle r^2 \rangle_{A'A} / \mu_{A'A}}$

electronic factors    nuclear factor

King linearity: eliminating the nuclear factor

$$\tilde{\nu}_{A'A}^{(2)} = K_{21} + F_{21} \tilde{\nu}_{A'A}^{(1)}$$

$$K_{21} := K_2 - F_{21} K_1, \quad F_{21} := F_2 / F_1$$

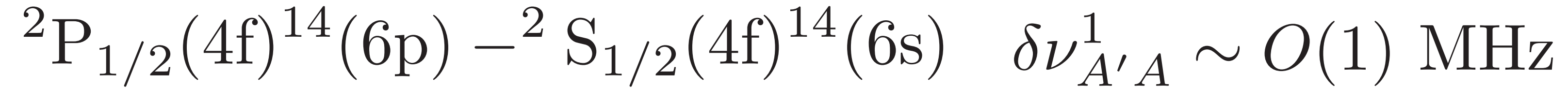
  $(\tilde{\nu}_{A'A}^{(1)}, \tilde{\nu}_{A'A}^{(2)})$  on a straight line, King plot

# Ex. Yb<sup>+</sup>

K. Mikami, MT, Y. Yamamoto EPJC77:896 (2017)

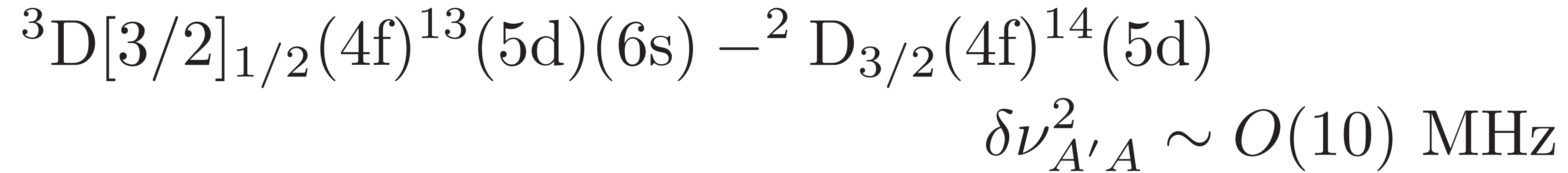
## Transition 1: 369 nm

Martensson-Pendrill et al. PRA49, 3351 (1994)



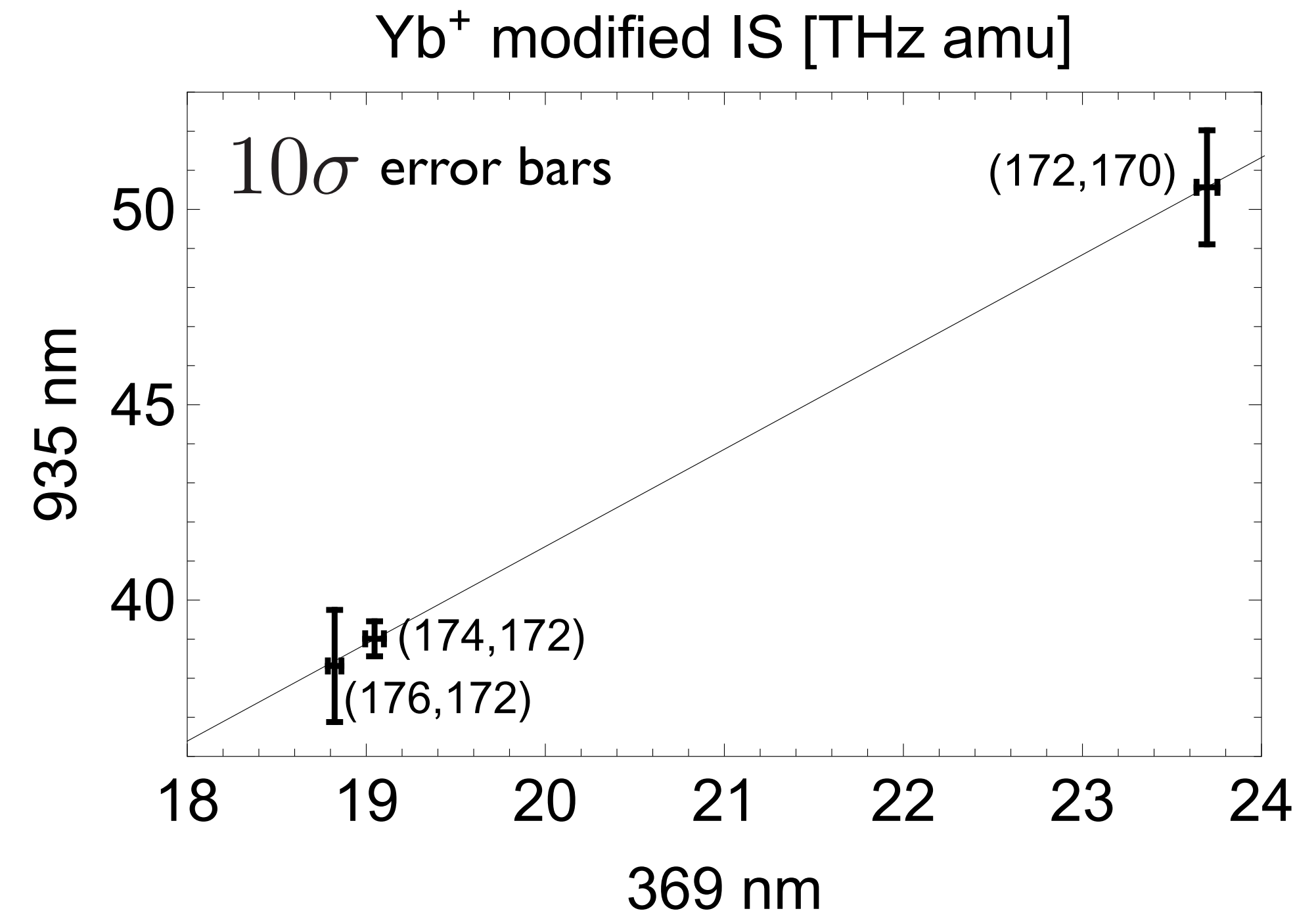
## Transition 2: 935 nm

Sugiyama et al. CPEM2000



## Isotope pairs

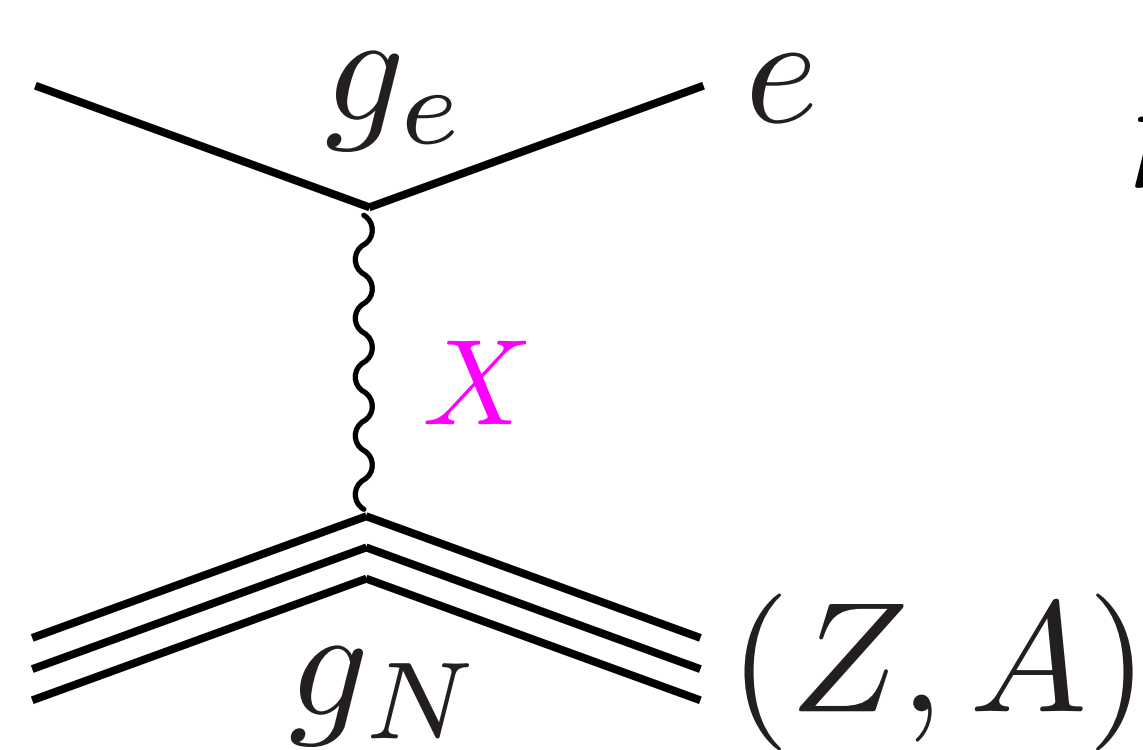
(172, 170), (174, 172), (176, 172)



# Nonlinearity

IS by new **neutron-electron interaction**

Delaunay et al. arXiv:1601.05087v2



$$\nu_{A'A}^{(t)} = \underbrace{K_t \mu_{A'A}}_{\text{MS}} + \underbrace{F_t \langle r^2 \rangle_{A'A}}_{\text{FS}} + \underbrace{X_t (A' - A)}_{\text{particle shift (PS)}}$$

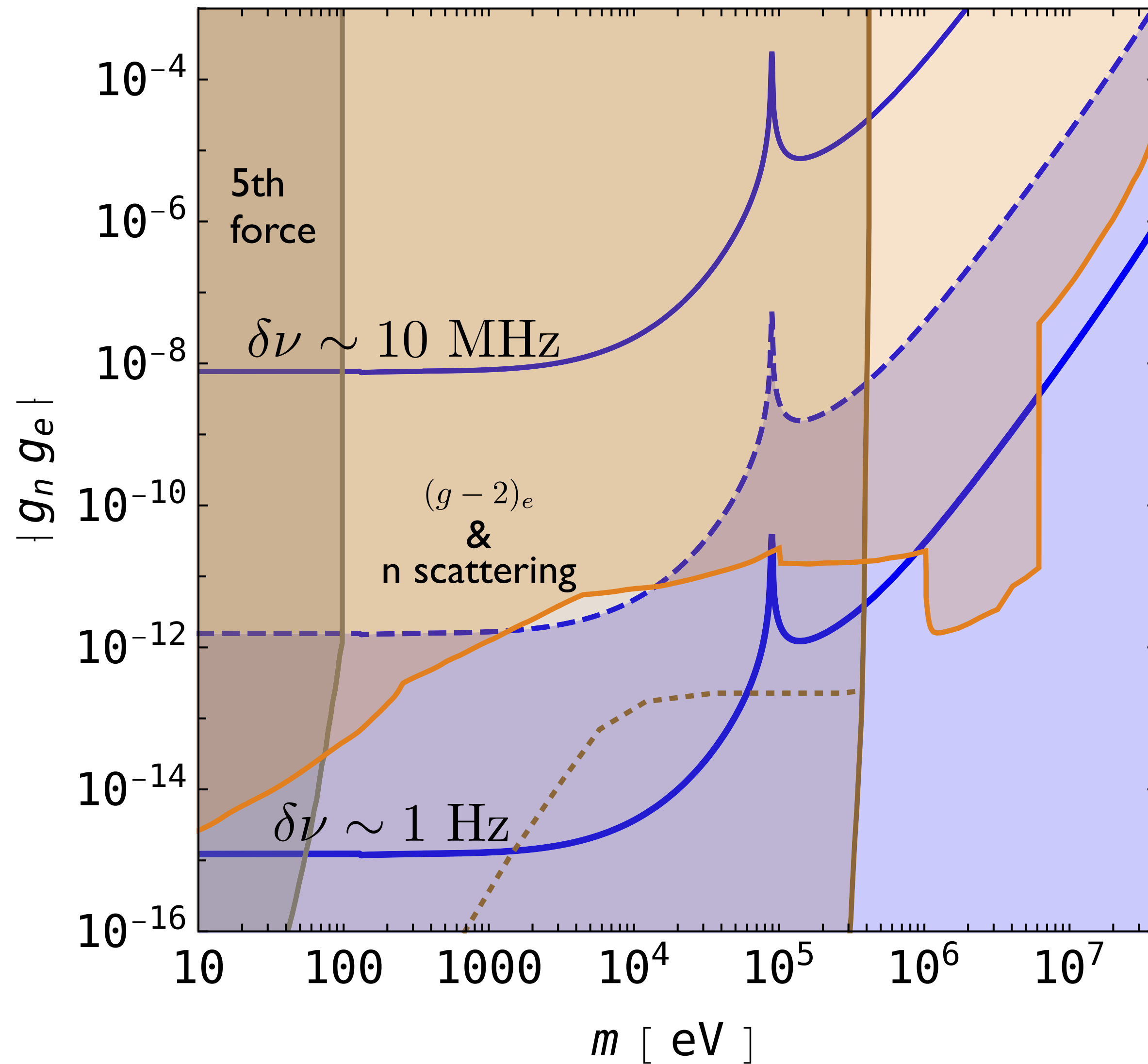
Nonlinearity due to **subleading FS**

$$\text{FS} = F_t \langle r^2 \rangle_{A'A} + \underbrace{F'_t [\langle r^2 \rangle_{A'A}]^2}_{\text{quadratic FS}} + \underbrace{G_t \langle r^4 \rangle_{A'A}}_{\text{higher moment}} + \dots$$

$$[\langle r^2 \rangle_{A'A}]^2 := (\langle r^2 \rangle_{A'A_0})^2 - (\langle r^2 \rangle_{AA_0})^2$$

# Ex. Yb<sup>+</sup>

MT, Y. Yamamoto PTEP I03B02 (2020)



**Transition 1: 369 nm** Martensson-Pendrill et al. PRA49, 3351 (1994)  
 $^2P_{1/2}(4f)^{14}(6p) - ^2S_{1/2}(4f)^{14}(6s)$   $\delta\nu_{A'A}^1 \sim O(1)$  MHz

**Transition 2: 935 nm** Sugiyama et al. CPEM2000  
 $^3D[3/2]_{1/2}(4f)^{13}(5d)(6s) - ^2D_{3/2}(4f)^{14}(5d)$   
 $\delta\nu_{A'A}^2 \sim O(10)$  MHz

— Yb<sup>+</sup> bounds  
 - - -  $\langle r^4 \rangle$  FS nonlinearity (SM BG)

**FSNL dominance:**

$$\delta\nu \lesssim 1 \text{ kHz}$$

**What about SM nonlinearity?**

**Precise calculation difficult**

# Generalized linearity

K. Mikami, MT, Y. Yamamoto EPJC77:896 (2017)

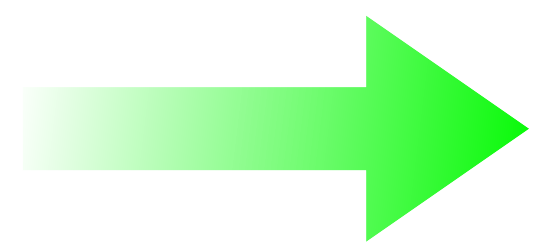
$$\nu_{A'A}^{(t)} = K_t \mu_{A'A} + F_t \langle r^2 \rangle_{A'A} + F'_t [\langle r^2 \rangle_{A'A}]^2 + X_t (A' - A)$$

3 transitions:  $t=1, 2, 3$

QFS

PS

$$\begin{pmatrix} \nu_{A'A}^{(1)} - X_1(A' - A) \\ \nu_{A'A}^{(2)} - X_2(A' - A) \\ \nu_{A'A}^{(3)} - X_3(A' - A) \end{pmatrix} = \begin{pmatrix} K_1 & F_1 & F'_1 \\ K_2 & F_2 & F'_2 \\ K_3 & F_3 & F'_3 \end{pmatrix} \begin{pmatrix} \mu_{A'A} \\ \langle r^2 \rangle_{A'A} \\ [\langle r^2 \rangle_{A'A}]^2 \end{pmatrix} =: M \begin{pmatrix} \mu_{A'A} \\ \langle r^2 \rangle_{A'A} \\ [\langle r^2 \rangle_{A'A}]^2 \end{pmatrix}$$



$$\begin{aligned} & (M^{-1})_{11} \nu_{A'A}^{(1)} + (M^{-1})_{12} \nu_{A'A}^{(2)} + (M^{-1})_{13} \nu_{A'A}^{(3)} \\ & - \{ (M^{-1})_{11} X_1 + (M^{-1})_{12} X_2 + (M^{-1})_{13} X_3 \} (A' - A) = \mu_{A'A} \end{aligned}$$

$$(\nu_{A'A}^{(1)}, \nu_{A'A}^{(2)}, \nu_{A'A}^{(3)}) / \mu_{A'A}$$

on a plane if  $X_t = 0$

$n$  transitions and  $n+1$  IS pairs  $\rightarrow$  NP search with  $n-2$  NL's removed

# Recent experiments: Yb<sup>+</sup> ion

Count et al. PRL 125, 123002 (2020)

Transition 1: 411 nm  
 $^2S_{1/2}(6s) - ^2D_{5/2}(5d)$

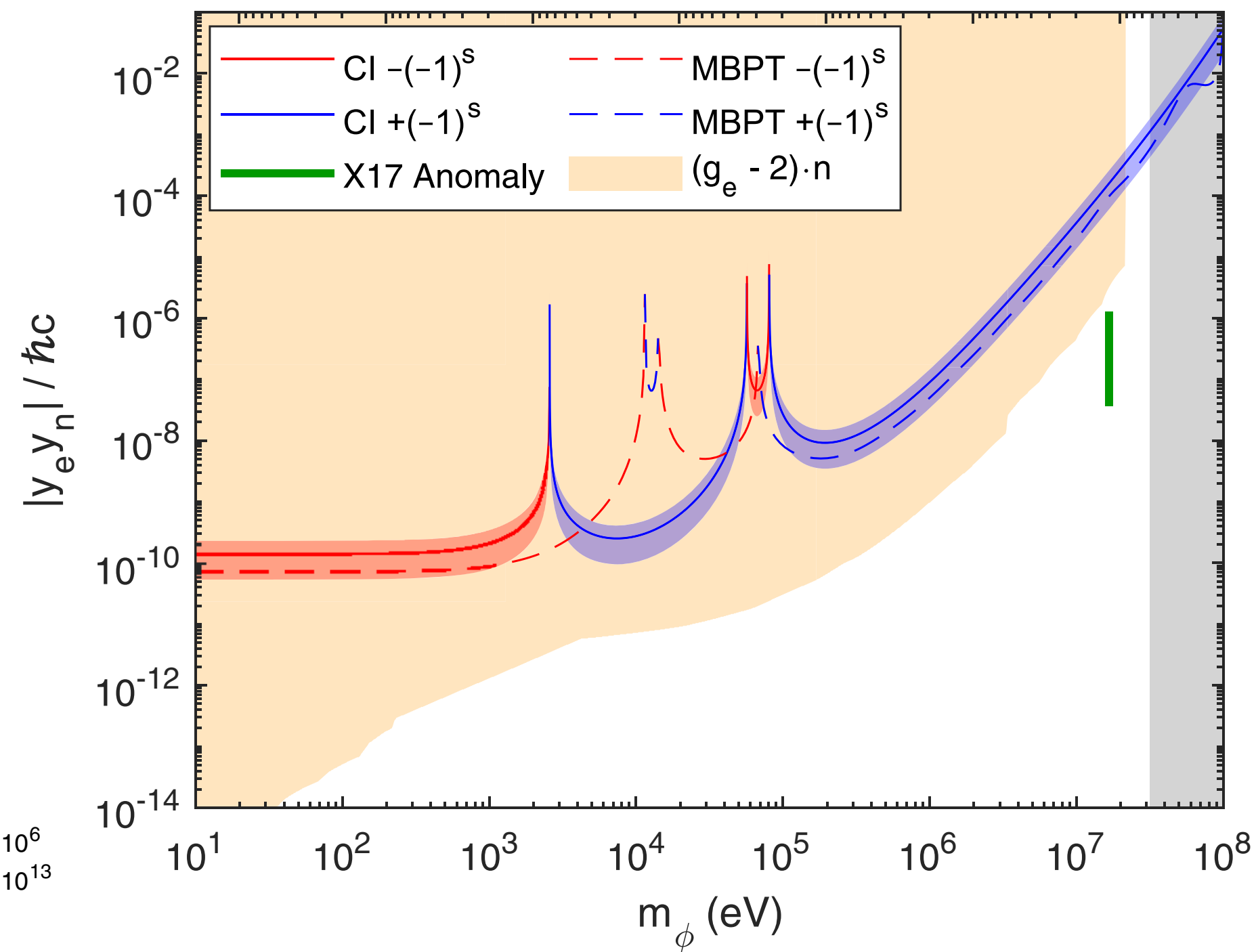
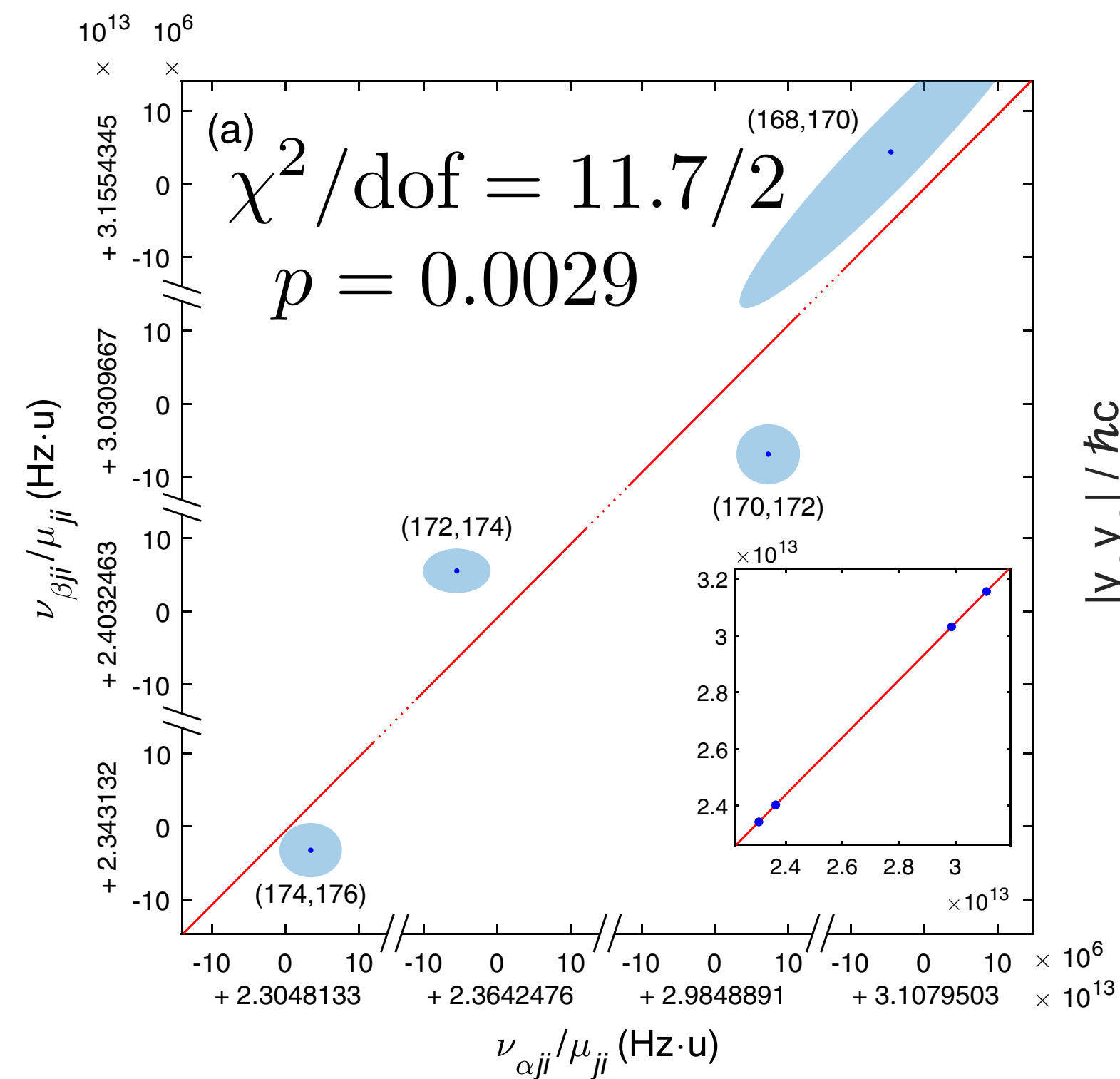
Transition 2: 436 nm  
 $^2S_{1/2}(6s) - ^2D_{3/2}(5d)$

4 indep. IS pairs  
**A=168,170,172,174,176**

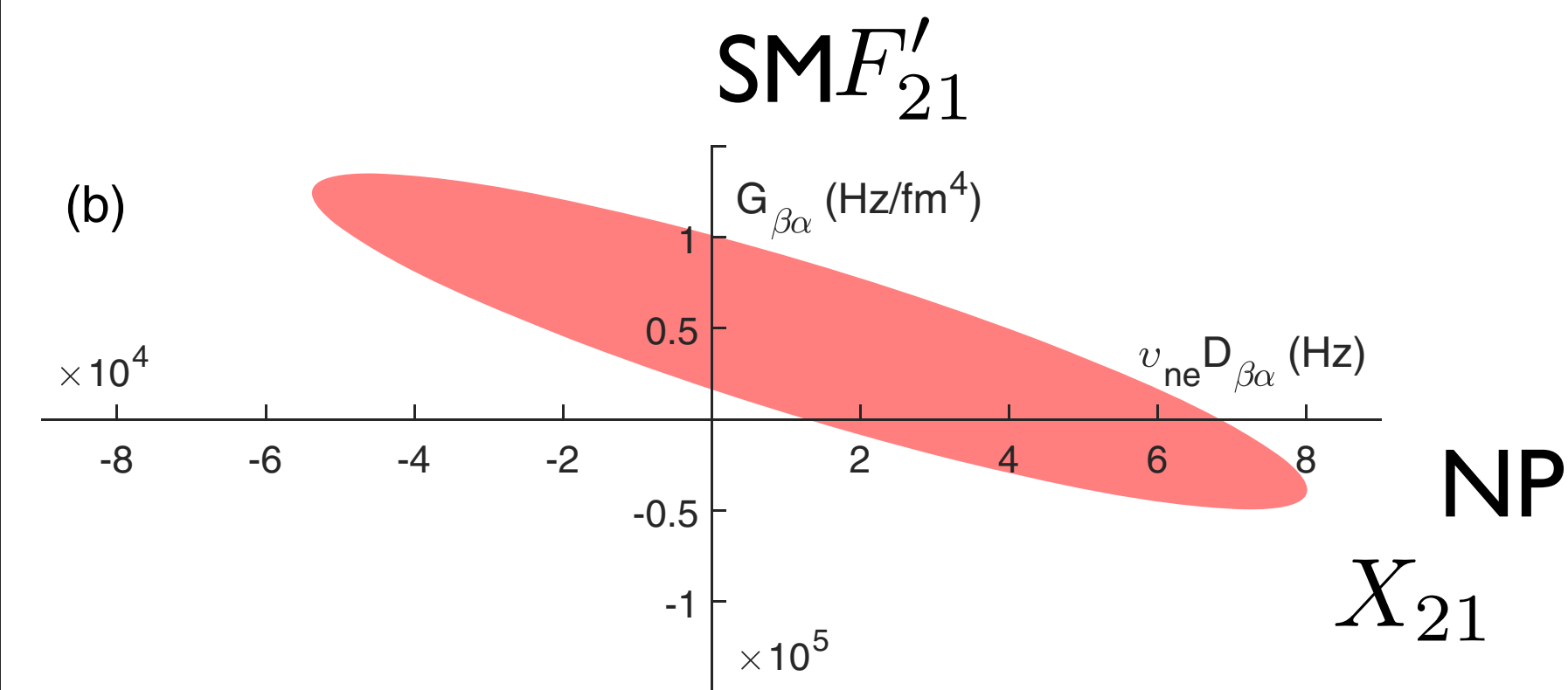
$\delta\nu \sim 300$  Hz

evidence for nonlinearity

new physics?

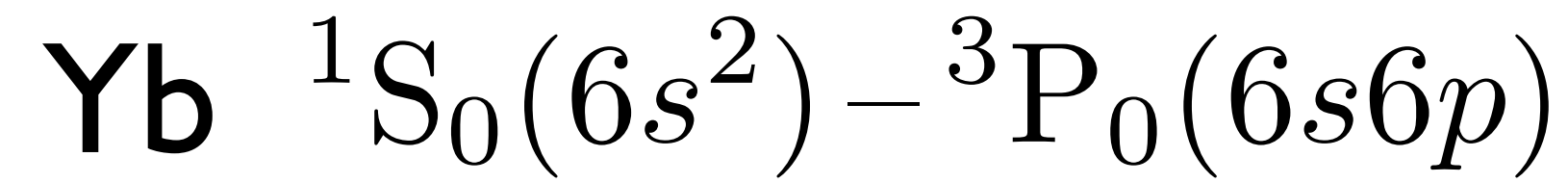


SM vs NP nonlinearities



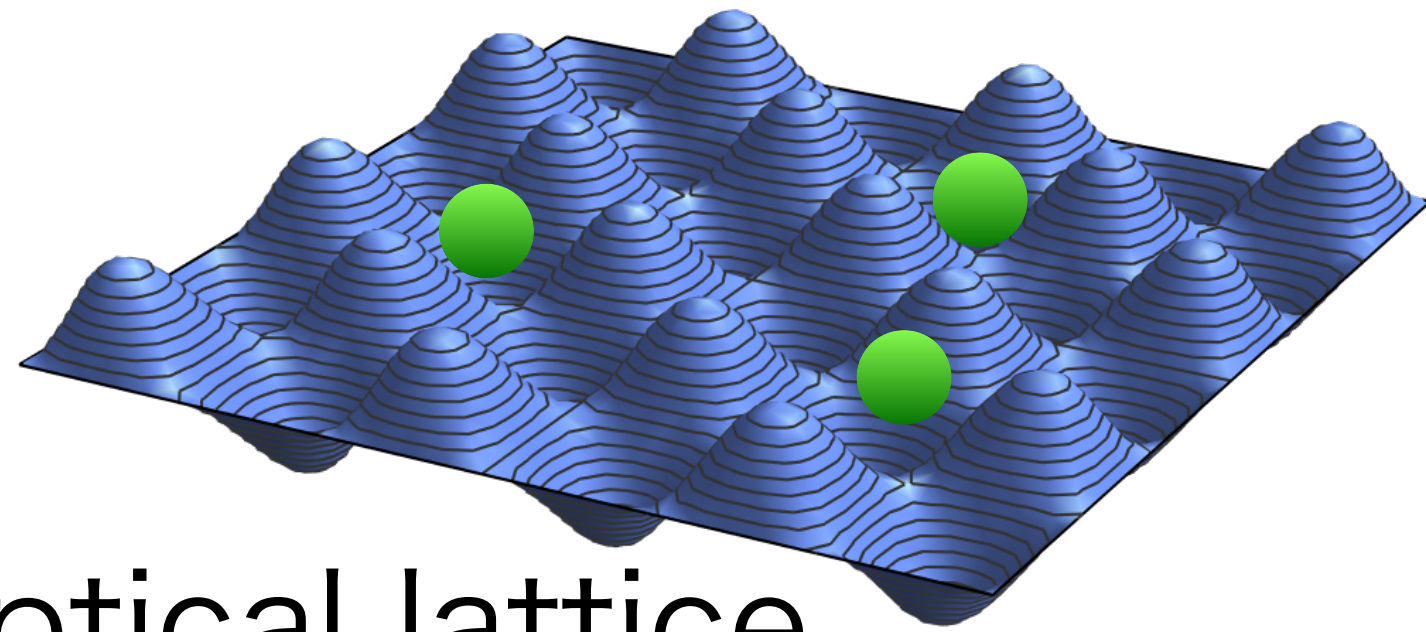


# Recent experiments: Yb atom K. Ono, MT et al. PRX 12, 021033 (2022)

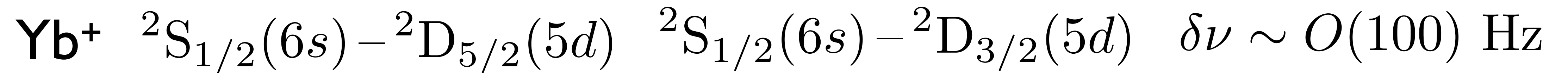
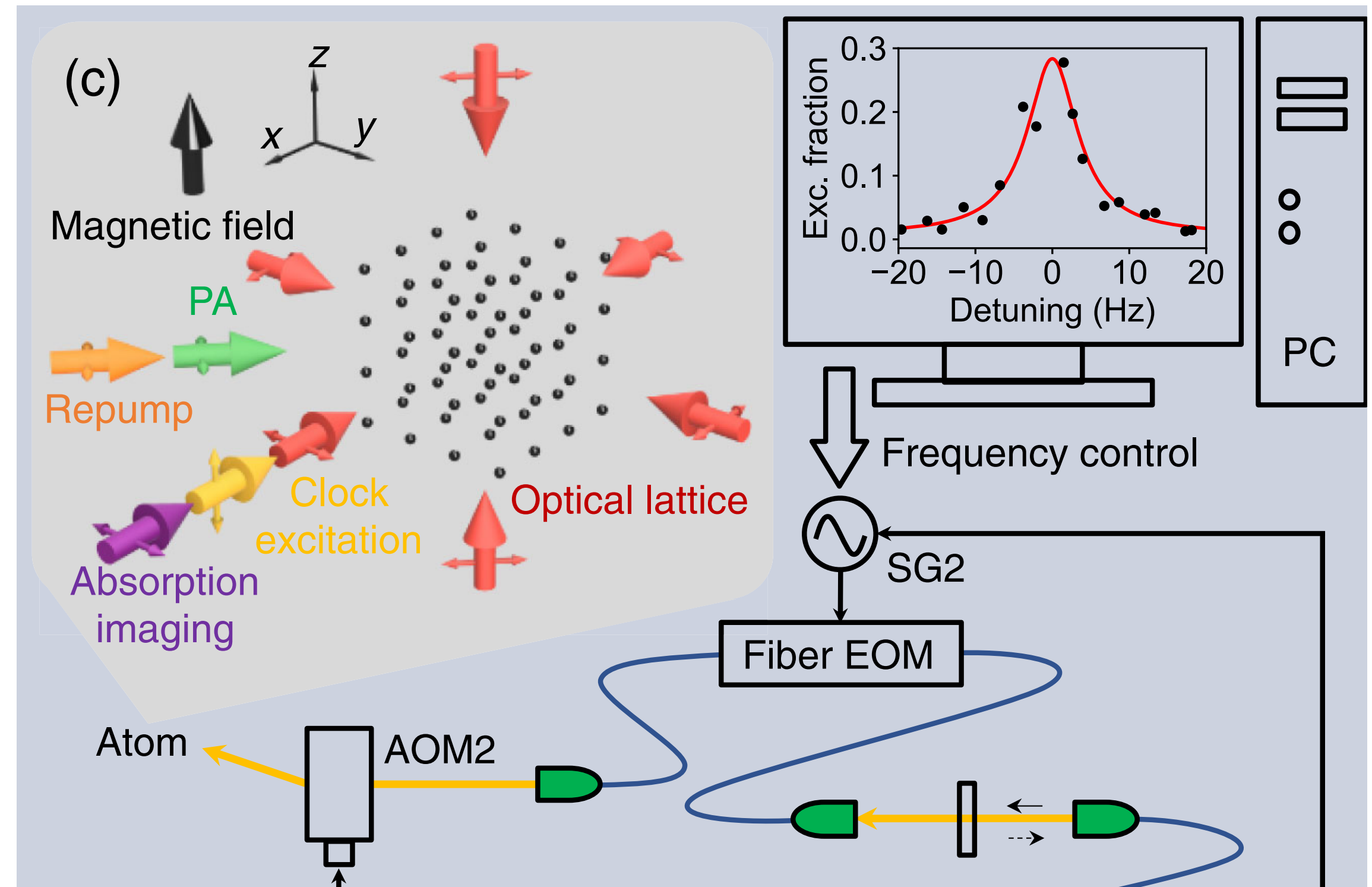


578 nm, 4 IS pairs

$\delta\nu \sim$  a few Hz



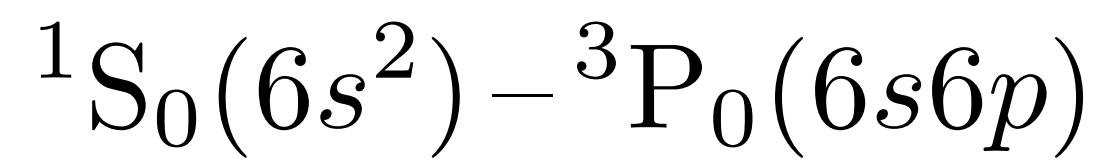
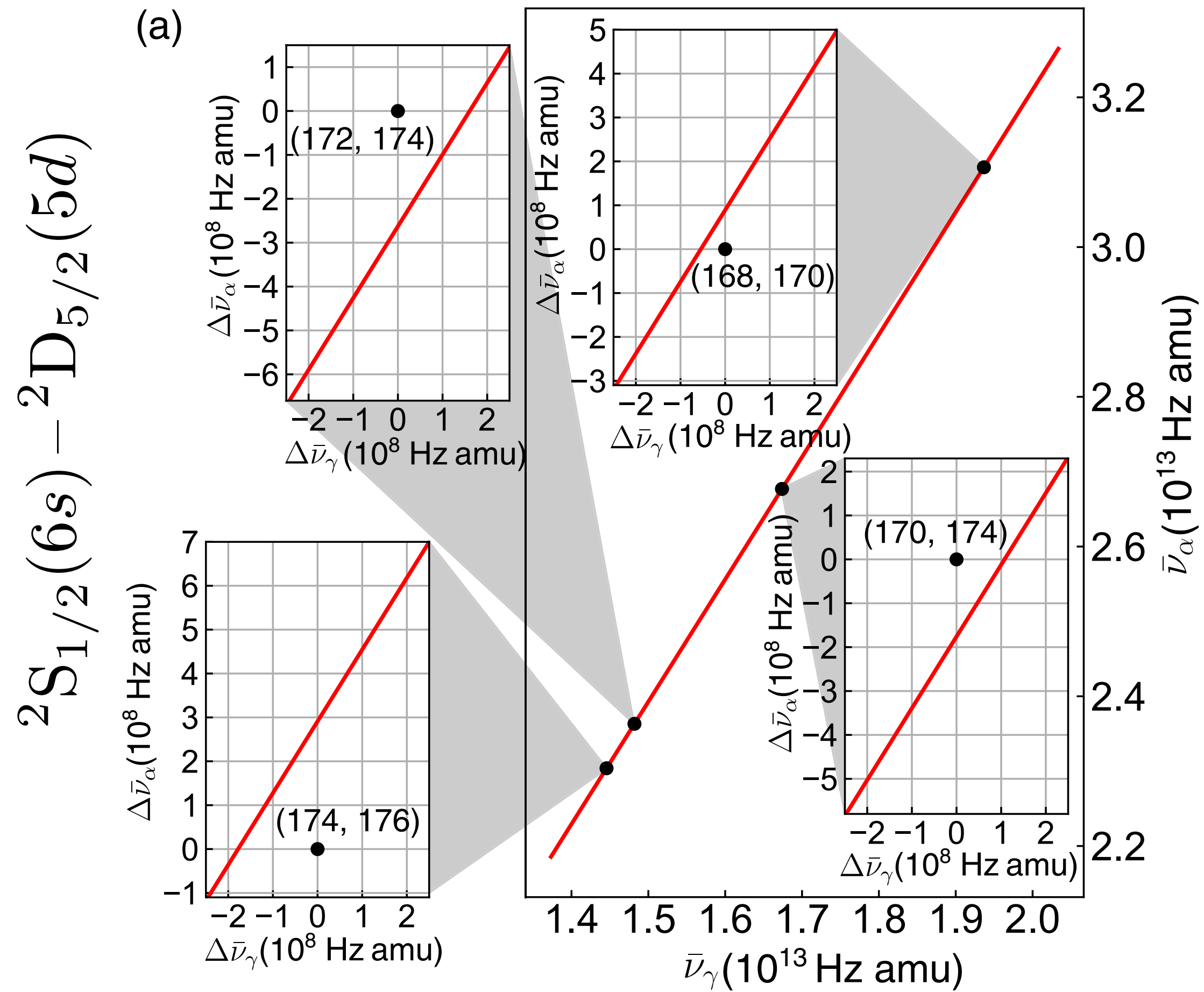
optical lattice



3 transitions, 4 IS pairs

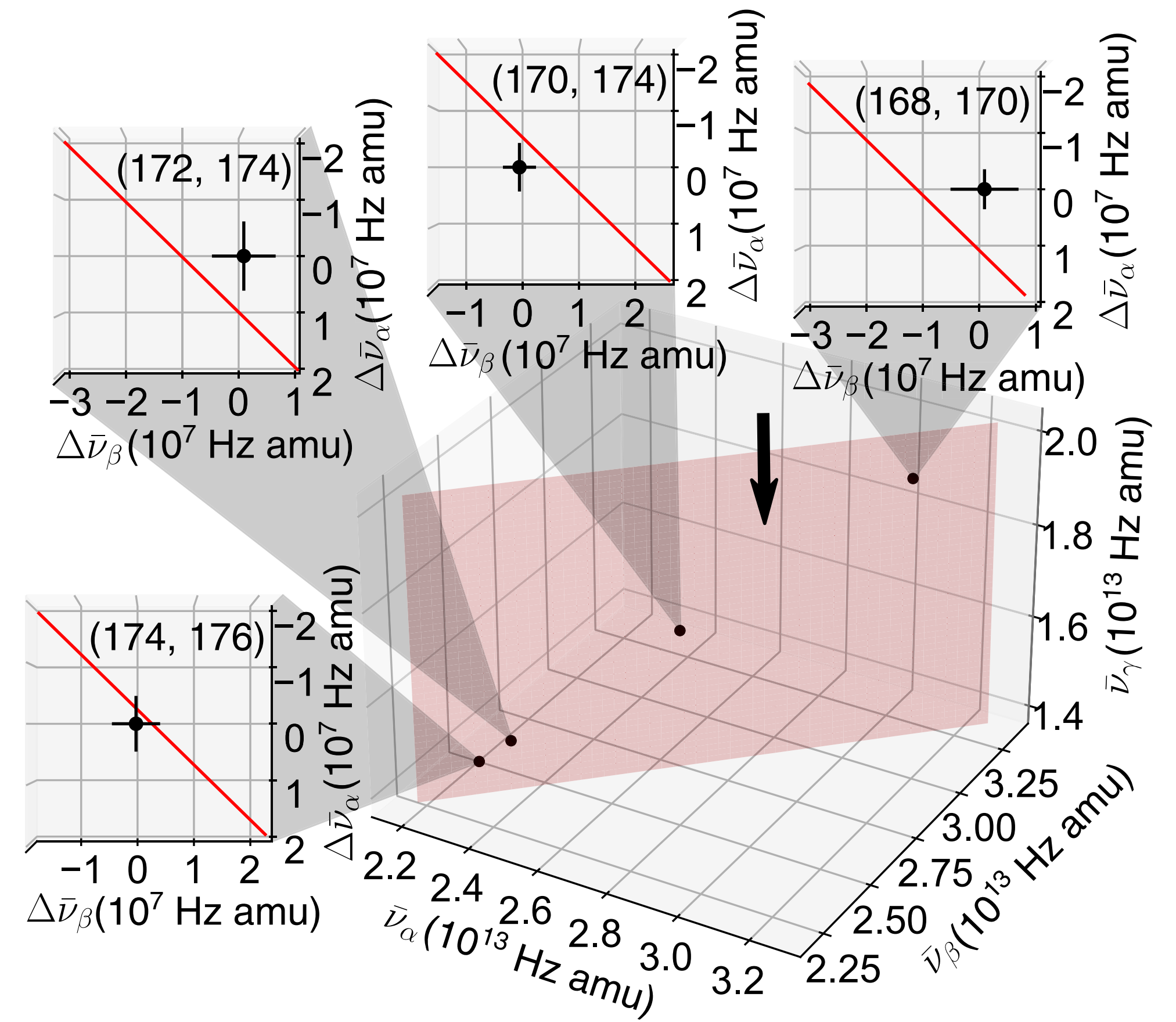
**➡ First new physics search with the generalized linearity**

# 2D analysis



$\chi^2/\text{dof} = 1.1 \times 10^4/3$

# 3D analysis



$\chi^2/\text{dof} = 15/3$

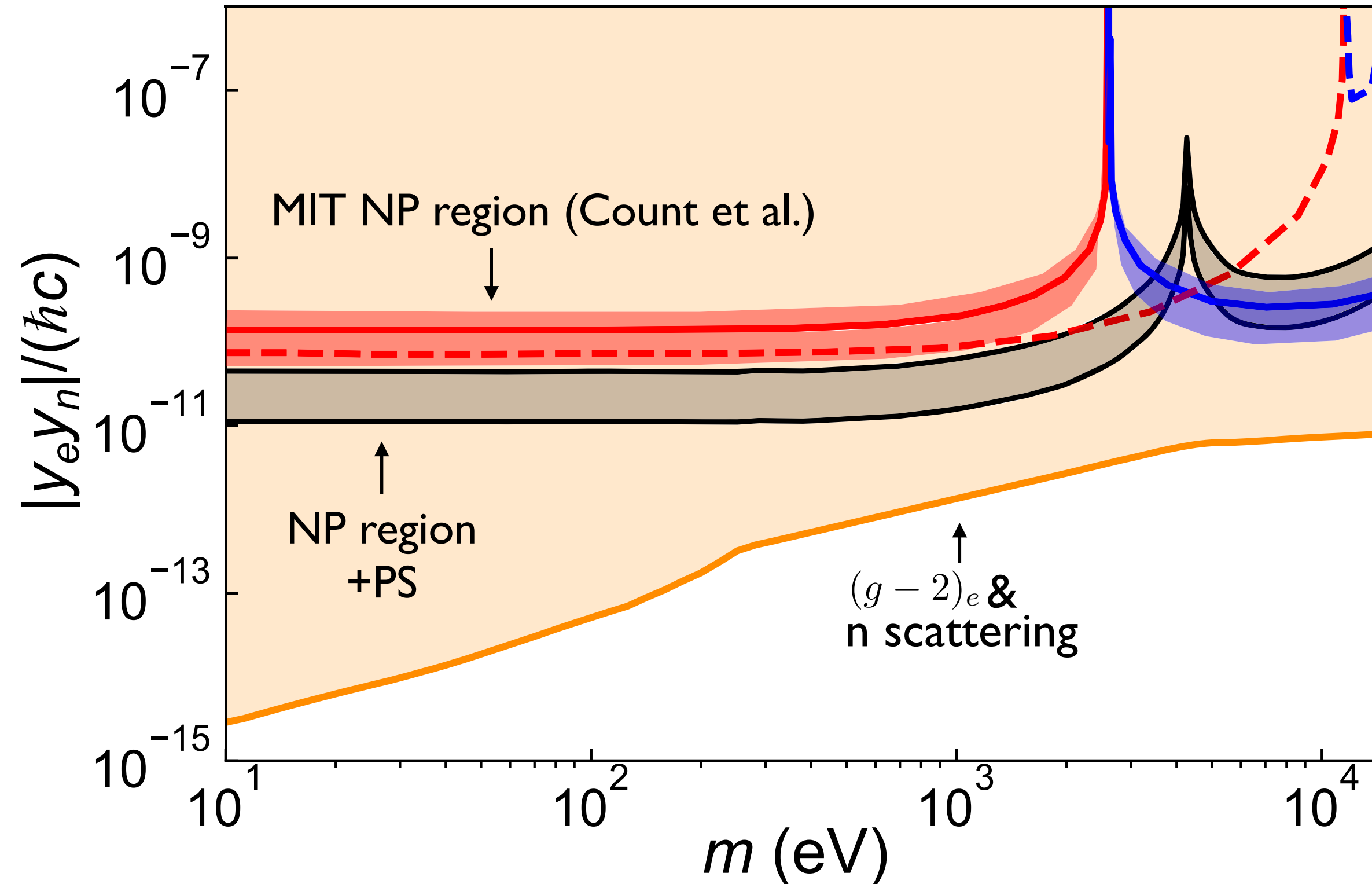


two or more NL sources

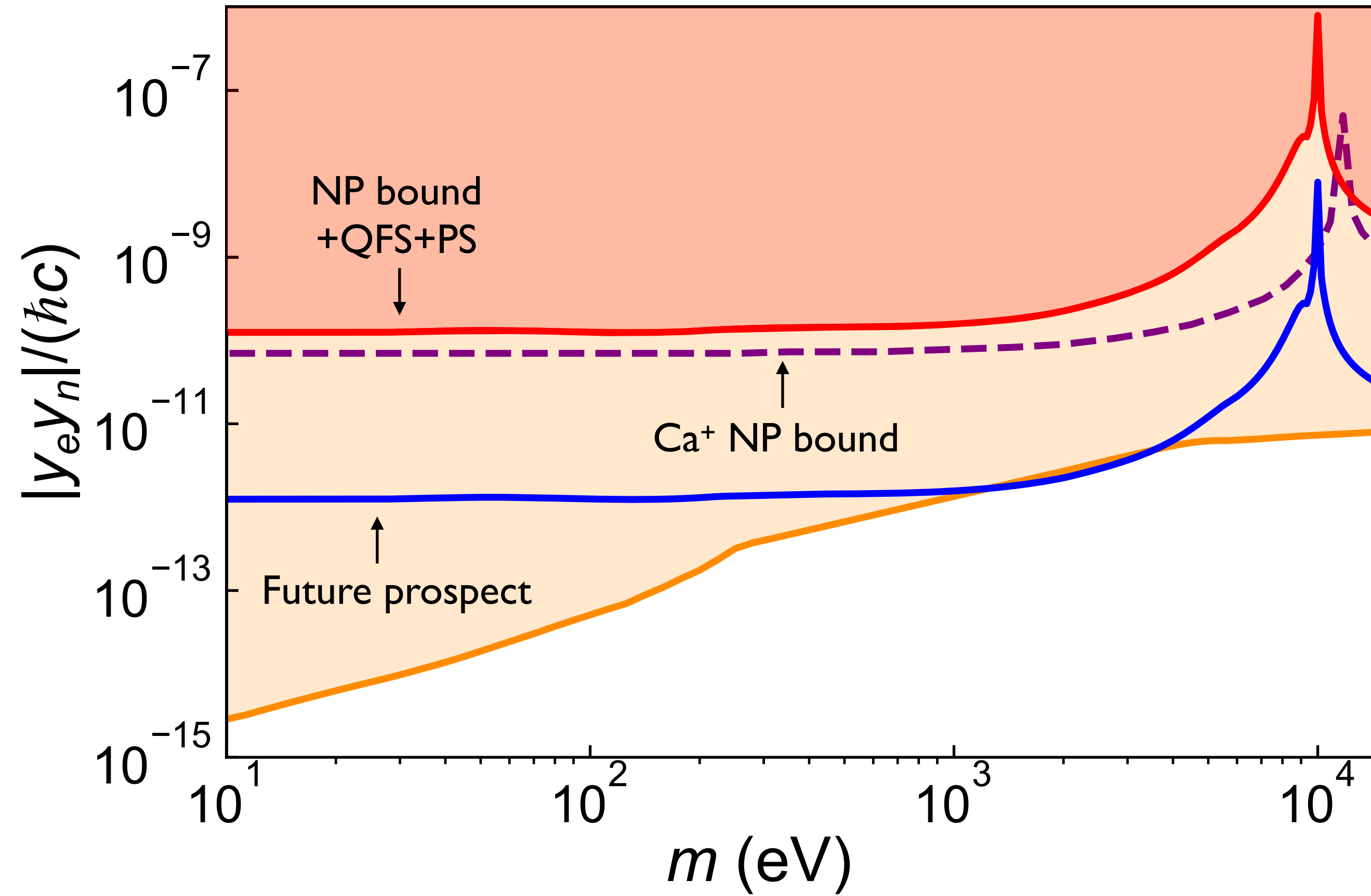
# Nonlinearity sources and new physics bound

One of NL sources is eliminated in 3D analysis.

+PS: Inconsistent with the existing constraints of PS



+QFS+PS  $\longrightarrow$  New physics bound





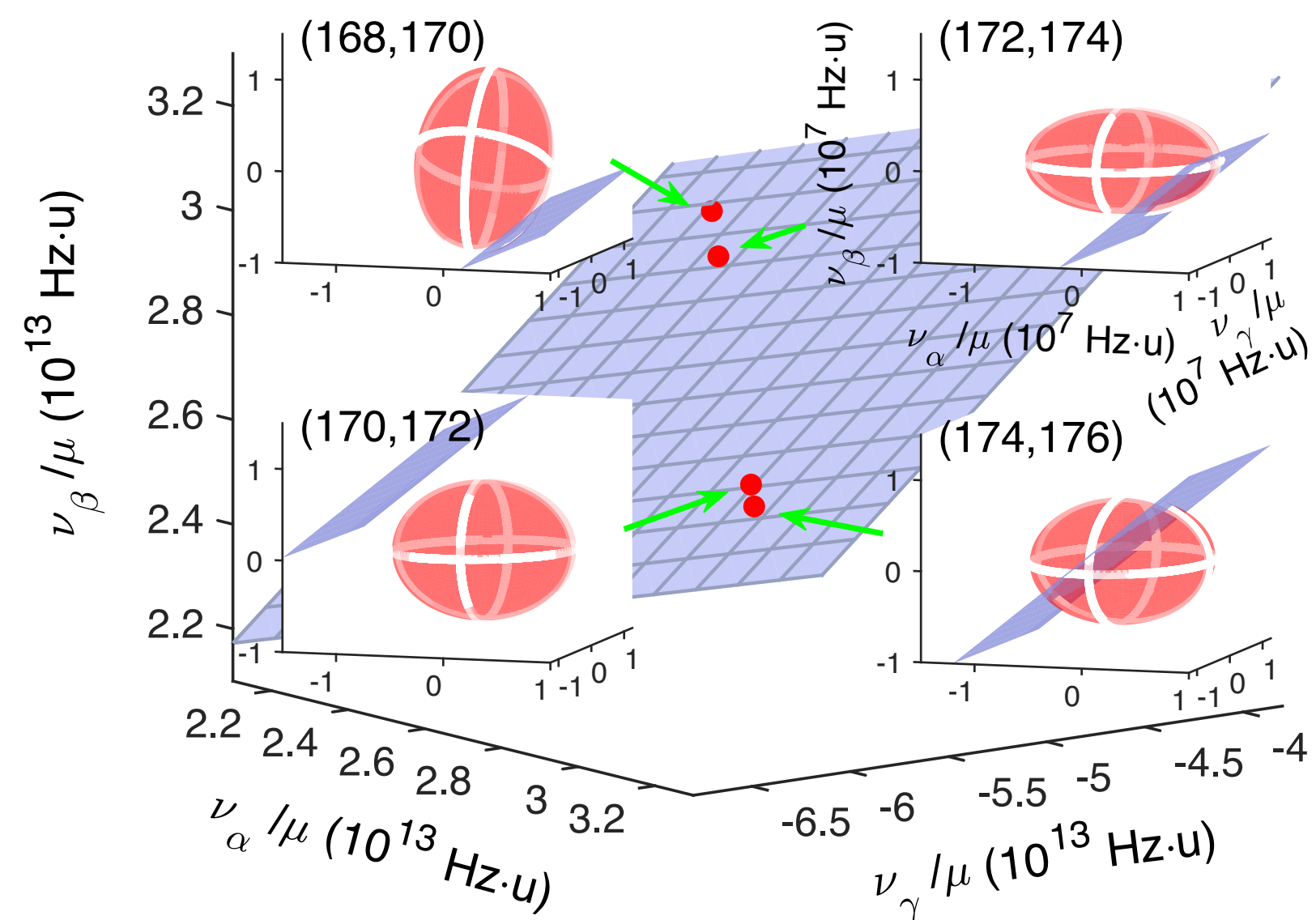
# Recent experiments (cont'd)

**Yb<sup>+</sup>** Hur et al. PRL 128, 163201 (2022)

$$^2S_{1/2}(4f)^{14}(6s) - ^2F_{7/2}(4f)^{13}(6s)^2$$

$$\delta\nu \sim 500 \text{ Hz}$$

3D analysis  $3.2 \sigma$

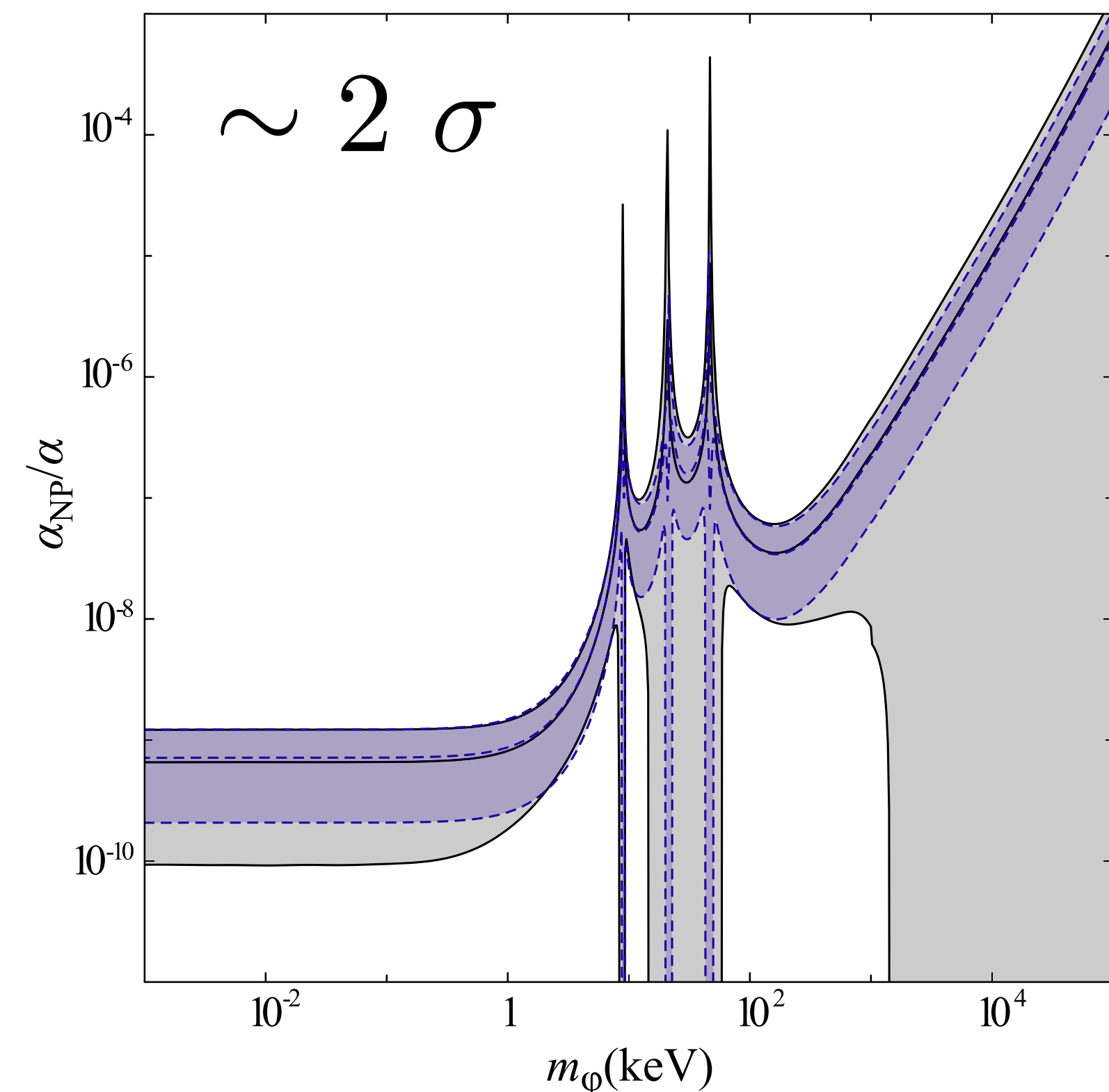


Consistent with our result

**Yb** Figueroa et al. PRL 128, 073001 (2022)

$$^1S_0(6s)^2 - ^1D_2(6s5d) \quad \delta\nu \sim O(100) \text{ Hz}$$

3D analysis: reduced significance



# Summary and outlook

- **Isotope shift and King linearity**  $\tilde{\nu}_{A'A}^{(2)} = K_{21} + F_{21}\tilde{\nu}_{A'A}^{(1)}$   
IS=MS+FS, linear relation of mIS of two transitions
- **Nonlinearities**: New physics and/or SM higher order
- **Generalized linearity**  
SM nonlinearity removed, improved sensitivity to new physics
- **Precise Yb IS measurements**  
Yb<sup>+</sup> ion O(100) Hz, Yb atom O(1) Hz  
New Yb O(1) Hz data expected
- **Data analysis**  
Unifying the 5 transitions of Yb, Yb<sup>+</sup> (ongoing)